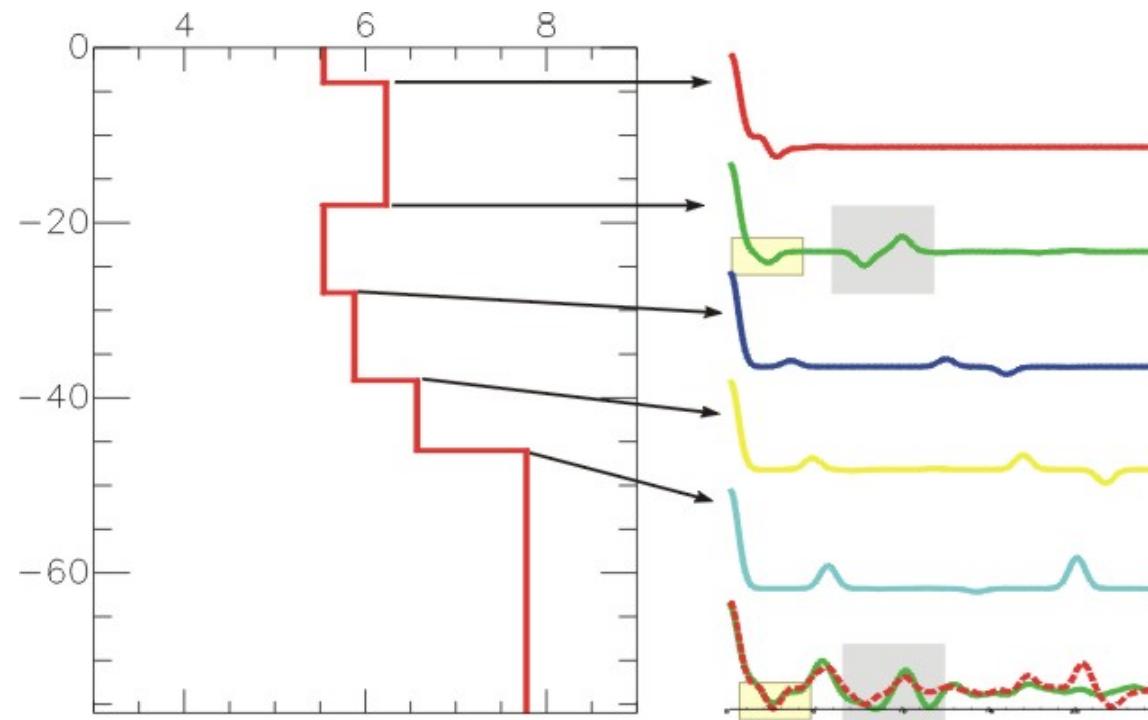
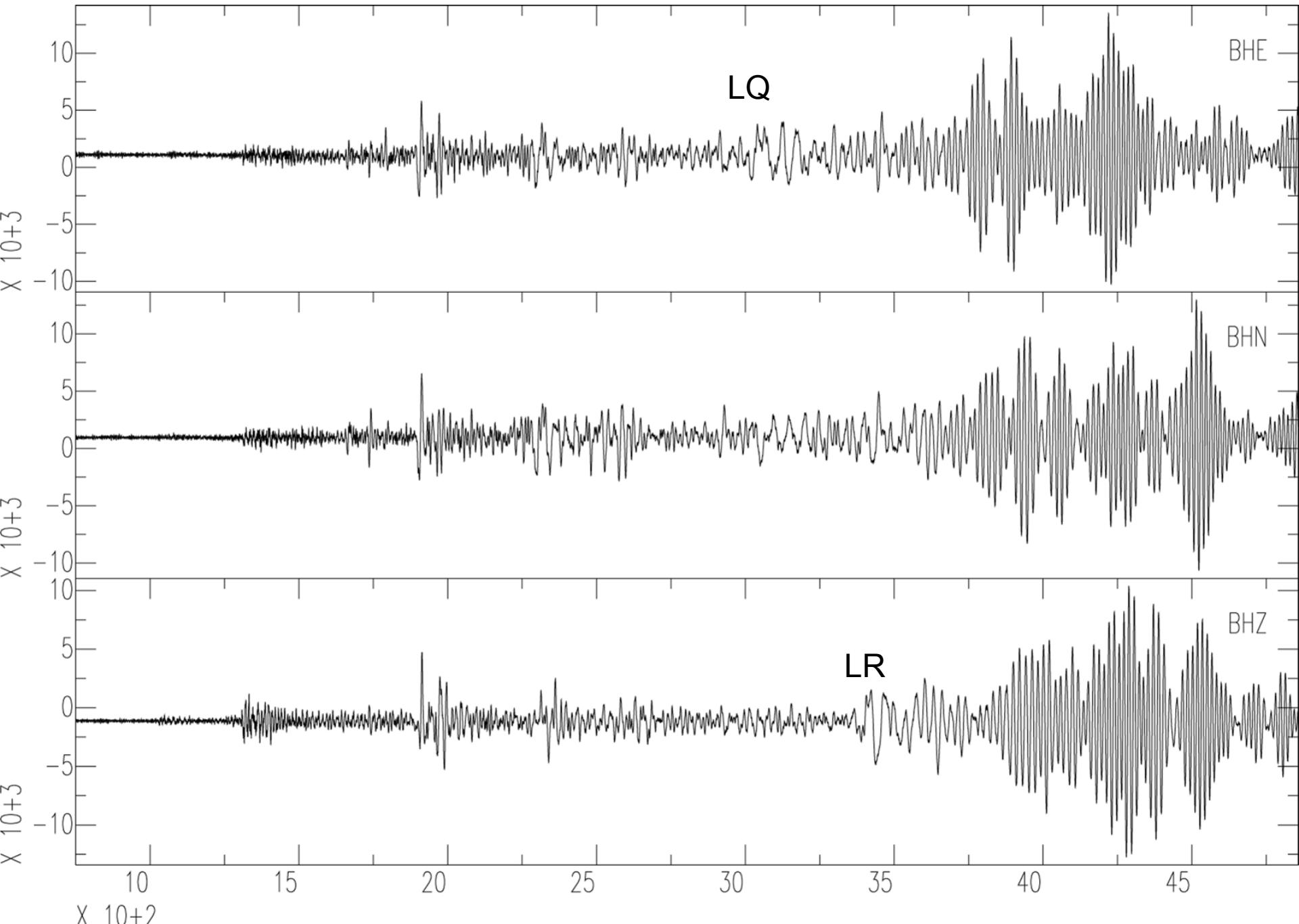


Introduction to Receiver Functions

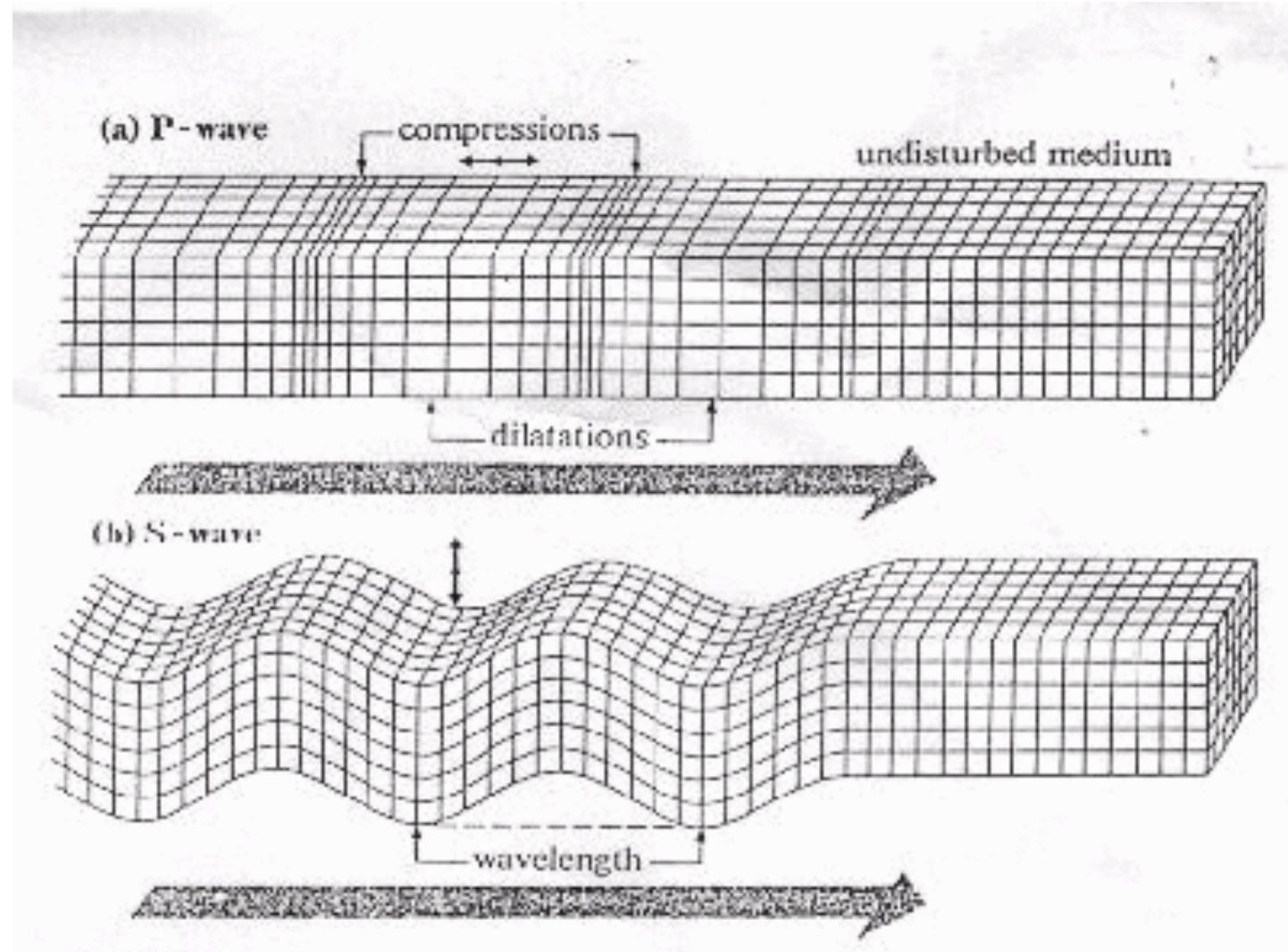
- I. Definition of a Receiver Function
- II. Slant Stacking
- III. Modeling of Receiver Functions
- IV. S-wave Receiver Functions





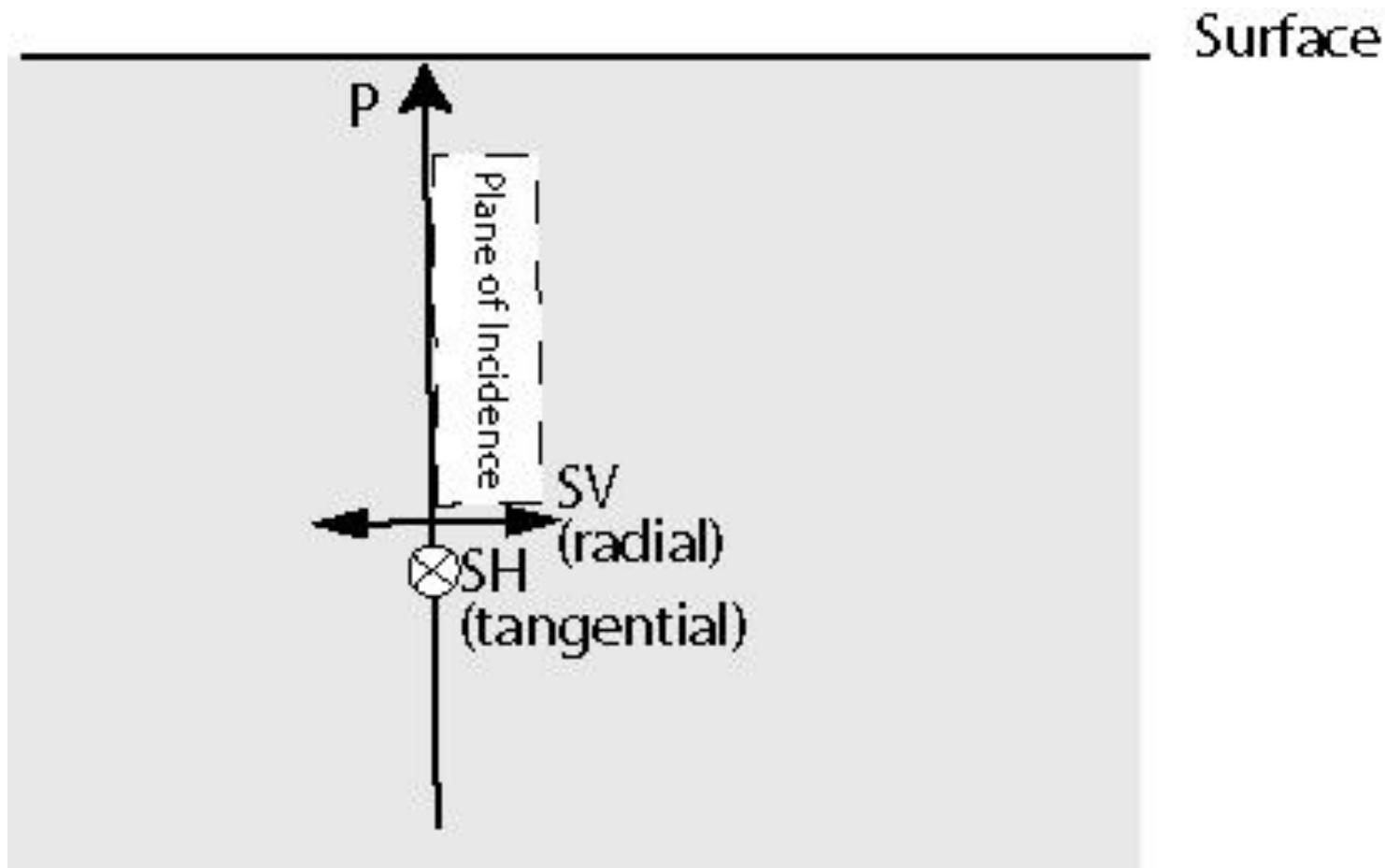
Body Waves

P and S waves Particle Motion

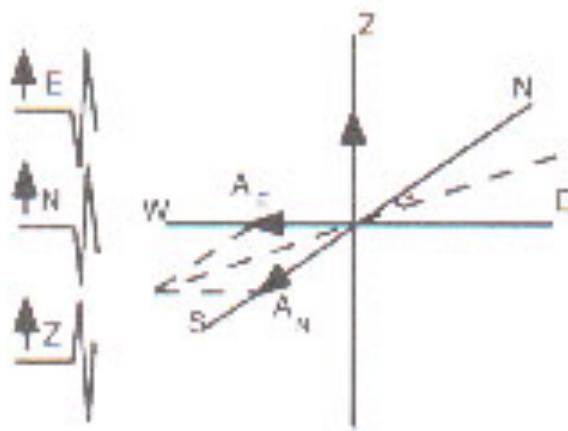
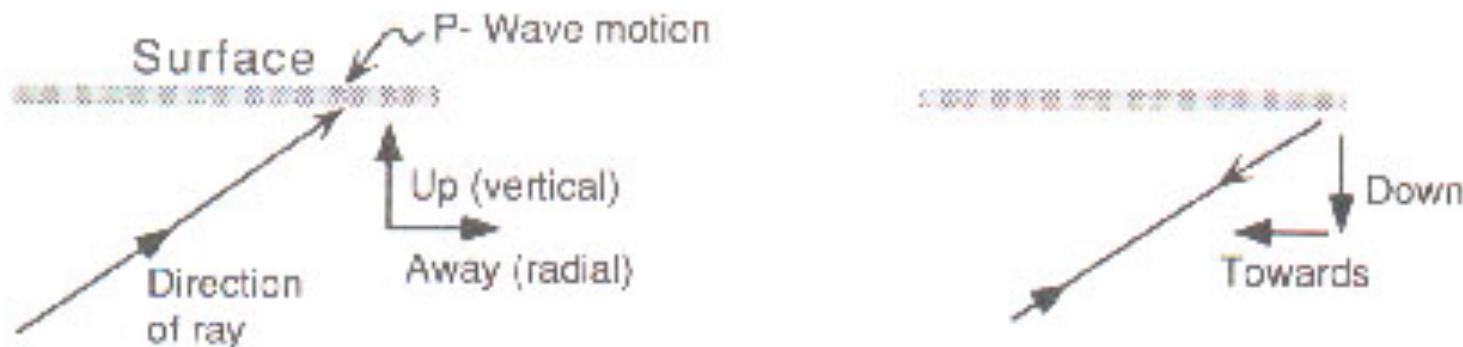


Body Waves: Polarization

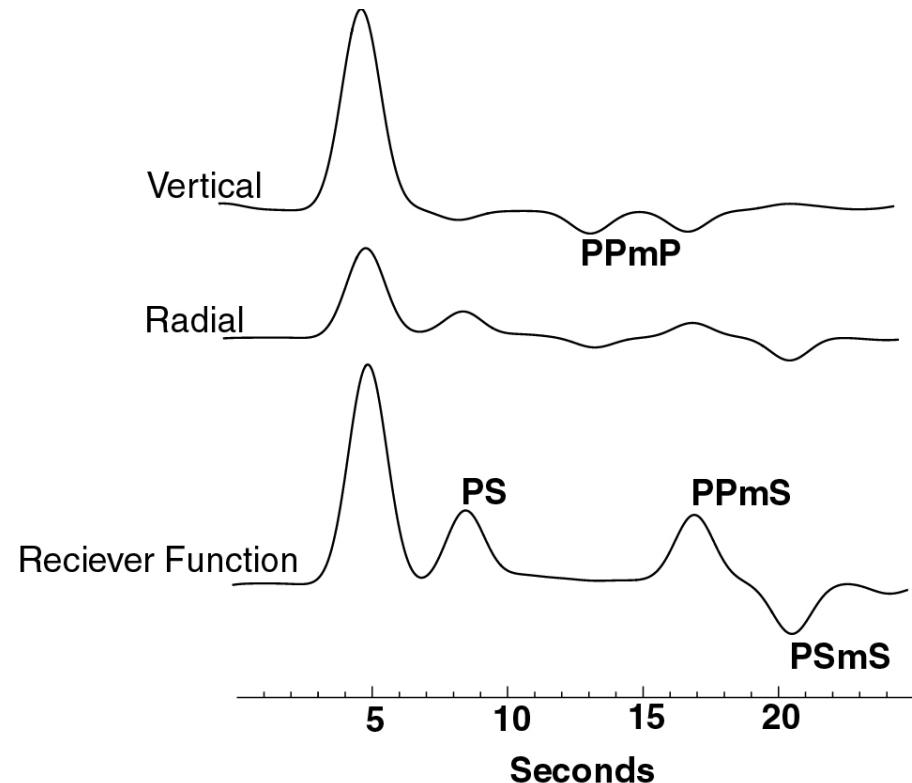
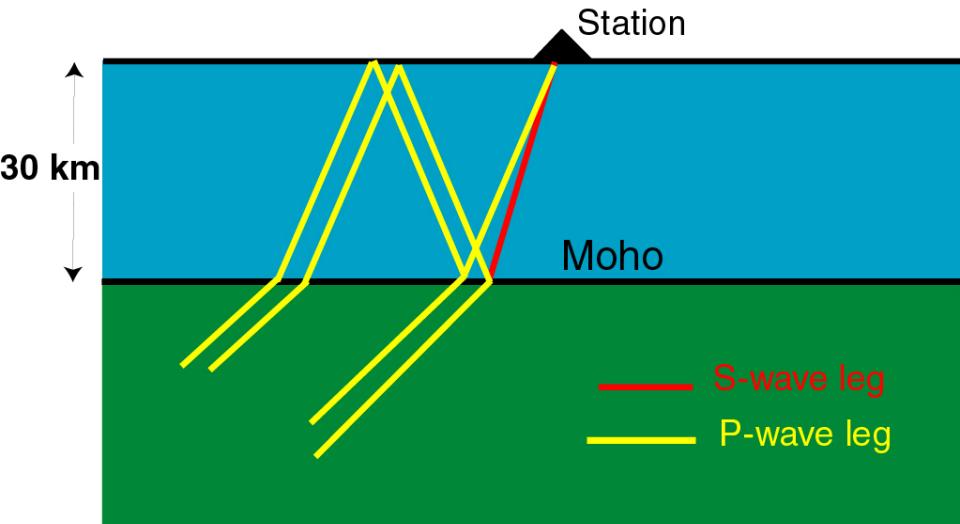
P, SV, and SH :



Seismic Body Waves



Receiver Functions



$$\begin{aligned}
 R(\omega) &= \frac{U_R(\omega)}{U_V(\omega)} = \frac{I(\omega) \cdot S(\omega) \cdot E_R^R(\omega) \cdot E_R^S(\omega)}{I(\omega) \cdot S(\omega) \cdot E_V^R(\omega) \cdot E_V^S(\omega)} \\
 &= \frac{E_R^R(\omega) \cdot E_R^S(\omega)}{E_V^R(\omega) \cdot E_V^S(\omega)} \cong \frac{E_R^R(\omega)}{E_V^R(\omega)}
 \end{aligned}$$

Fundamental Assumptions

- Plane Wave Approximation
- The PS is primarily recorded on the radial
(the vertical component is negligible)
- These assumptions lead to the result that
the receiver function ***only*** includes P-to-S
converted energy

$$\hat{z}_k = \frac{z_k}{z_0}, \hat{r}_k = \frac{r_k}{r_0}$$

$$E_R(\omega) = r_0 \left[1 + \hat{r}_p e^{-i\omega t_p} + \hat{r}_{ps} e^{-i\omega t_{ps}} \right]$$

$$E_Z(\omega) = z_0 \left[1 + \hat{z}_p e^{-i\omega t_p} + \hat{z}_{ps} e^{-i\omega t_{ps}} \right]$$

we can assume $z_{ps} \ll 1$

$$R(\omega) = \frac{r_0}{z_0} \frac{1 + \hat{r}_p e^{-i\omega t_p} + \hat{r}_{ps} e^{-i\omega t_{ps}}}{1 + \hat{z}_p e^{-i\omega t_p}}$$

using the binomial expansion : $(1+x)^{-1} = 1 - x + x^2 + \dots$

$$R(\omega) \cong \frac{r_0}{z_0} \left(1 - \hat{z}_p e^{-i\omega t_p} \right) \left(1 + \hat{r}_p e^{-i\omega t_p} + \hat{r}_{ps} e^{-i\omega t_{ps}} \right)$$

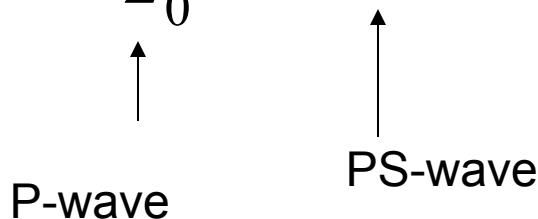
$$= \frac{r_0}{z_0} \left(1 - \hat{z}_p e^{-i\omega t_p} - \hat{z}_p \hat{r}_p e^{-2i\omega t_p} - \hat{z}_p \hat{r}_{ps} e^{-i\omega(t_p+t_{ps})} + \hat{r}_p e^{-i\omega t_p} + \hat{r}_{ps} e^{-i\omega t_{ps}} \right)$$

neglecting higher order terms

$$= \frac{r_0}{z_0} \left(1 - \hat{z}_p e^{-i\omega t_p} + \hat{r}_p e^{-i\omega t_p} + \hat{r}_{ps} e^{-i\omega t_{ps}} \right)$$

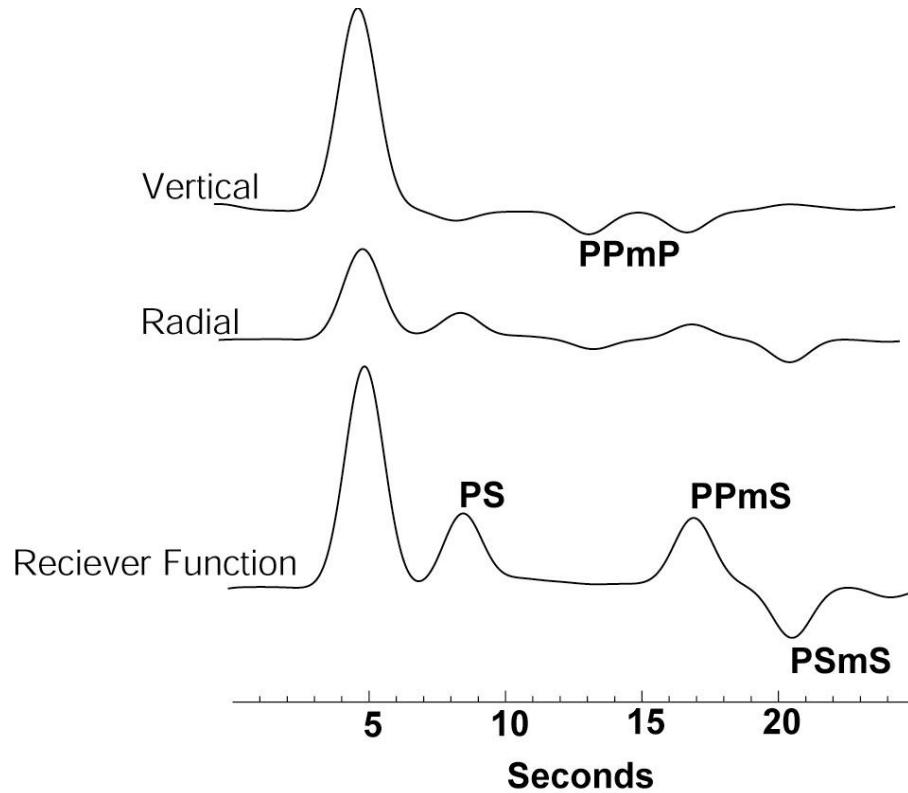
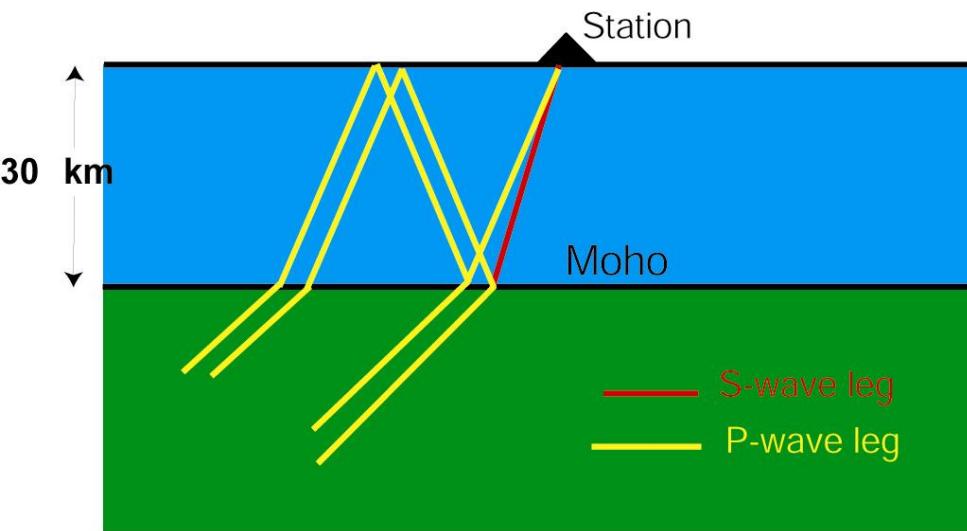
plane wave means $\hat{z}_p = \hat{r}_p$

$$\therefore R(\omega) = \frac{r_0}{z_0} + r_{ps} e^{-i\omega t_{ps}}$$



What factors influence r_0/z_0 ?

Isolating Mode Conversions



Methods of Deconvoltuion

- Spectral Division: water level technique
(Burdick and Langston, 1978)
- Stacked Time Domain Deconvolution
(Baker et al., 1996)
- Individual Iterative Time Domain
Deconvolution
(Liggoria and Ammon, 1999)

WATER-LEVEL DECONVOLUTION

- > water-level deconvolution introduced by Clayton & Wiggins

$$R(t) = \mathcal{F}^{-1} \left\{ \frac{P^*(\omega)S(\omega)}{\max(P^*(\omega)P(\omega), cP_{\max}^*P_{\max})} \right\}$$

- > for small c approaches deconvolution, large c approaches scaled cross-correlation
- > similar to damped least squares solution

$$R(t) = \mathcal{F}^{-1} \left\{ \frac{P^*(\omega)S(\omega)}{P^*(\omega)P(\omega) + \delta} \right\}$$

Time Domain CONVOLUTIONAL MODEL

$$u_{in}(\mathbf{x}, t) = S(t) \otimes G_{in}(\mathbf{x}, t; \mathbf{p}_\perp)$$

$$u_{in}(\mathbf{x}, t)$$

- observed displacement seismogram

$$S(t)$$

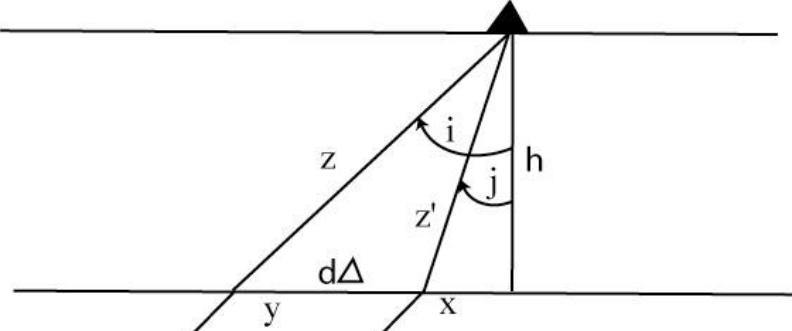
- effective source (includes source-side scattering)

$$G_{in}(\mathbf{x}, t; \mathbf{p}_\perp)$$

- Green's function (receiver side response to an impulsive plane wave with horizontal slowness)

$$\mathbf{p}_\perp$$

PS-phase Move-out:



$$t_{PS-P} = t_{PS} - t_P = \frac{x}{V_s} - \frac{y}{V_p} - dT = \frac{h}{V_s \cos j} - \frac{h}{V_p \cos i} - dT$$

using

$$\sin i = V_p p, \sin j = V_s p \quad \therefore \cos i = \sqrt{1 - p^2 V_p^2}, \cos j = \sqrt{1 - p^2 V_s^2},$$

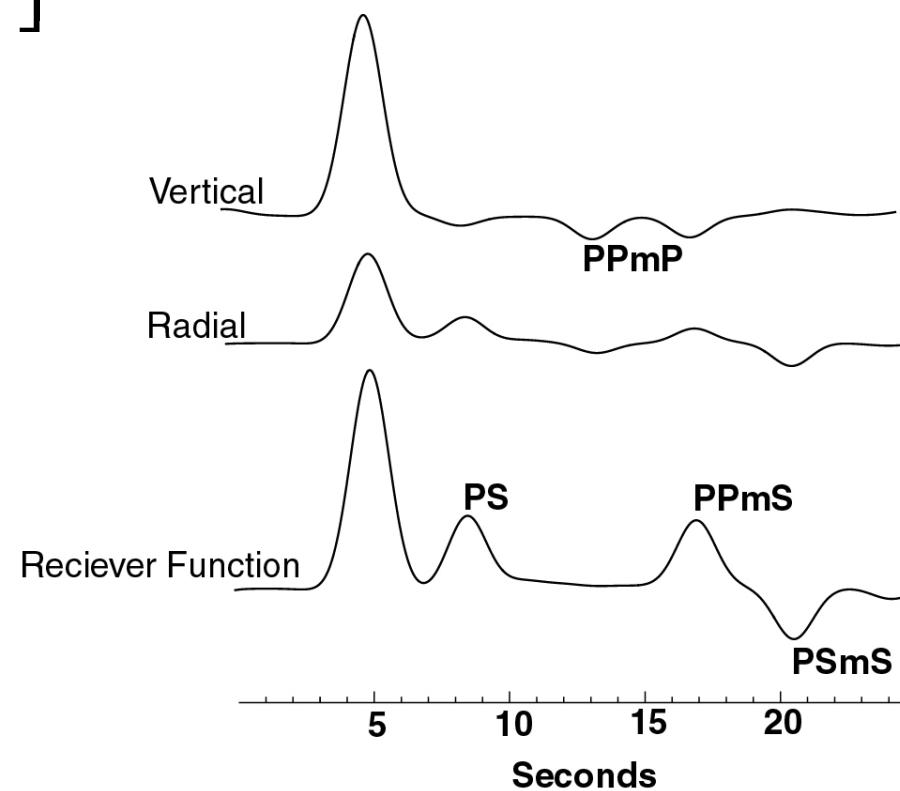
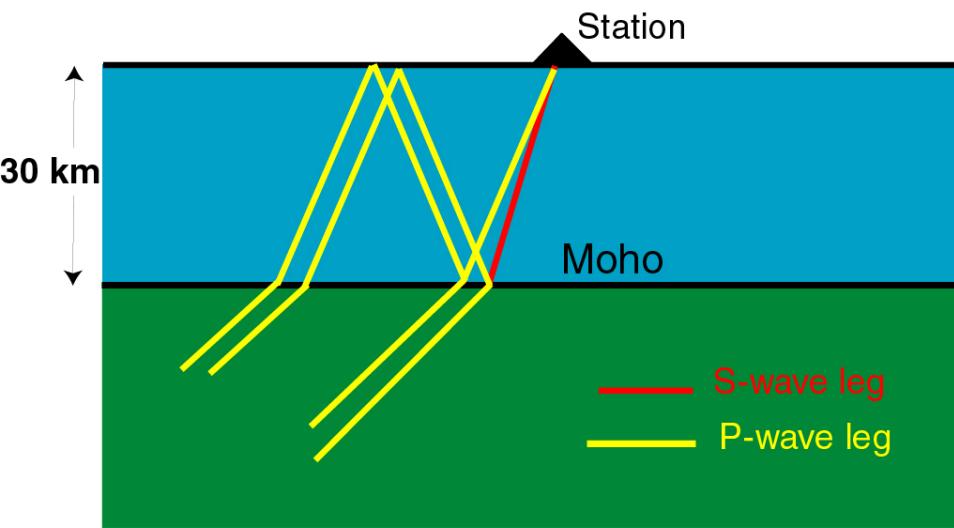
$$\text{then } t_{PS-P} = \frac{h\sqrt{1 - p^2 V_s^2}}{V_s} - \frac{h\sqrt{1 - p^2 V_p^2}}{V_p} - dT$$

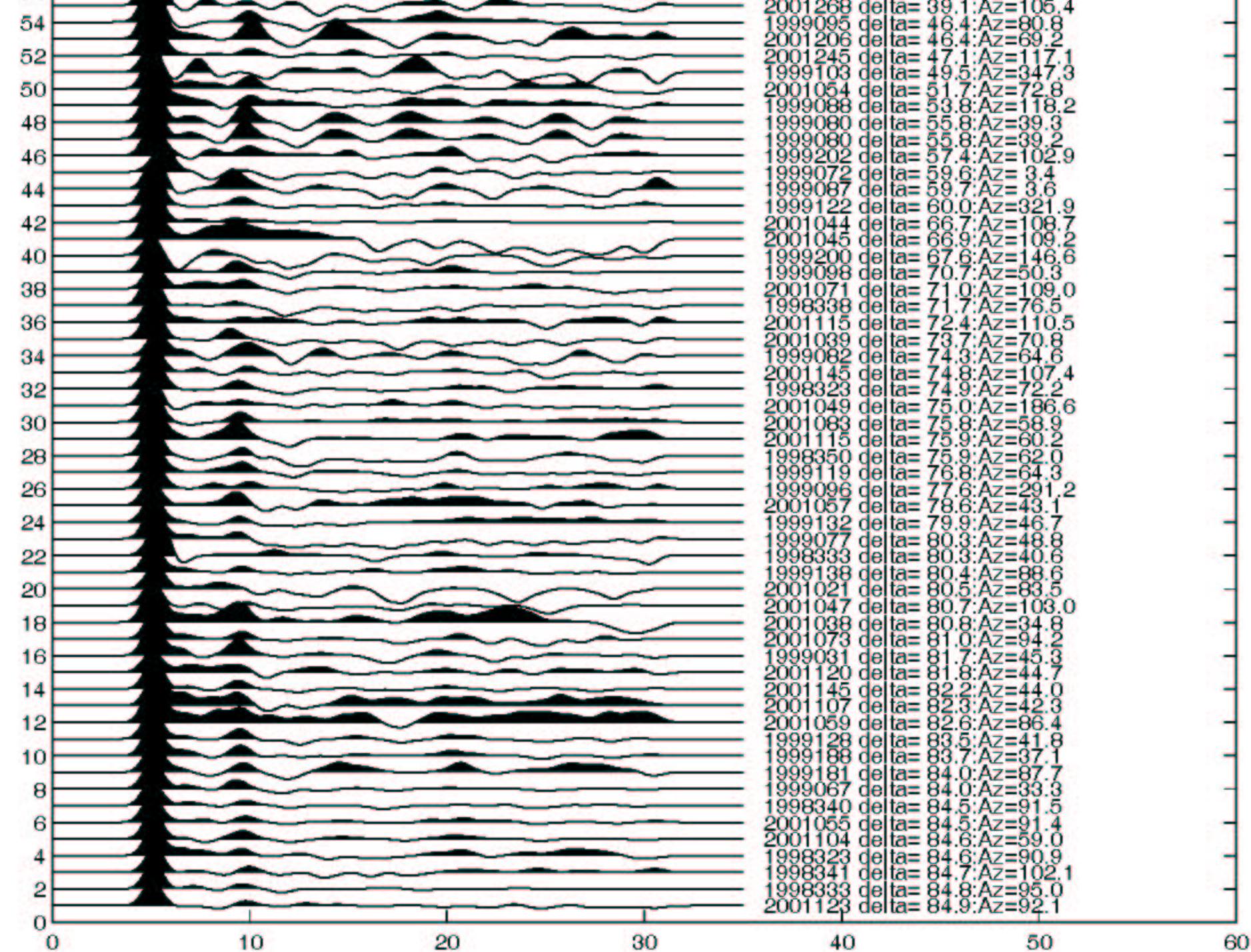
$$= \frac{h \left(V_p \sqrt{1 - p^2 V_s^2} - V_s \sqrt{1 - p^2 V_p^2} \right)}{V_s V_p} - pd\Delta$$

Crustal Multiples:

$$t_{PPS} = h \left[\sqrt{V_s^{-2} - p^2} + \sqrt{V_p^{-2} - p^2} \right]$$

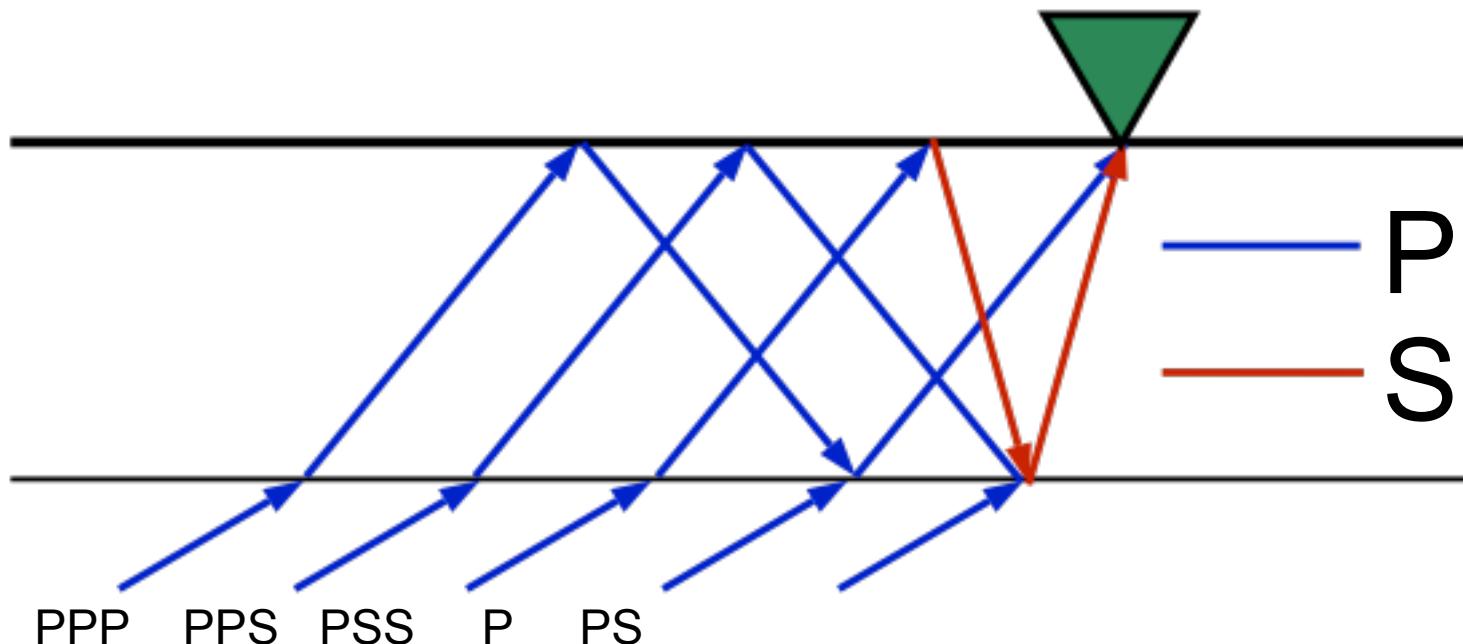
$$t_{PSS} = t_{PPS} + 2h \left[\sqrt{V_s^{-2} - p^2} \right]$$



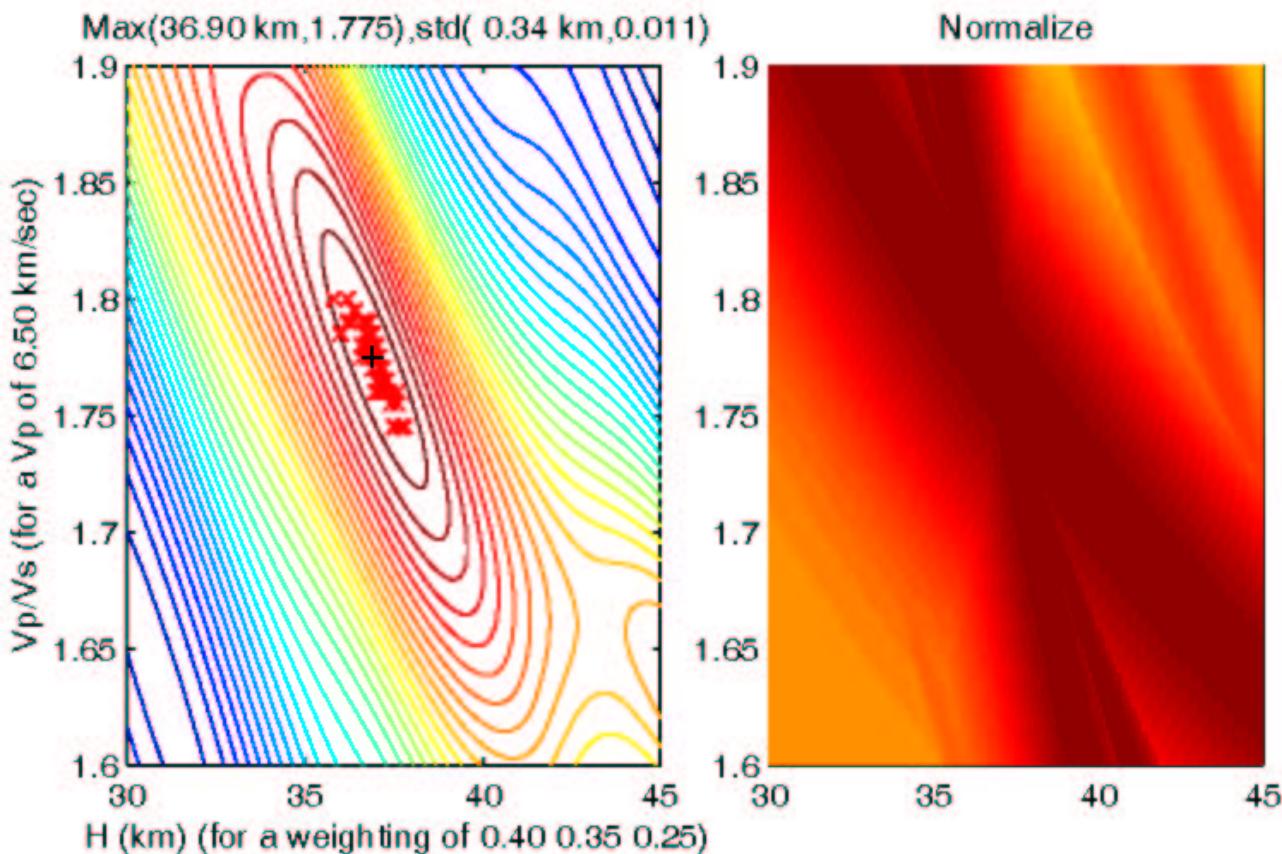


SCATTERING GEOMETRY

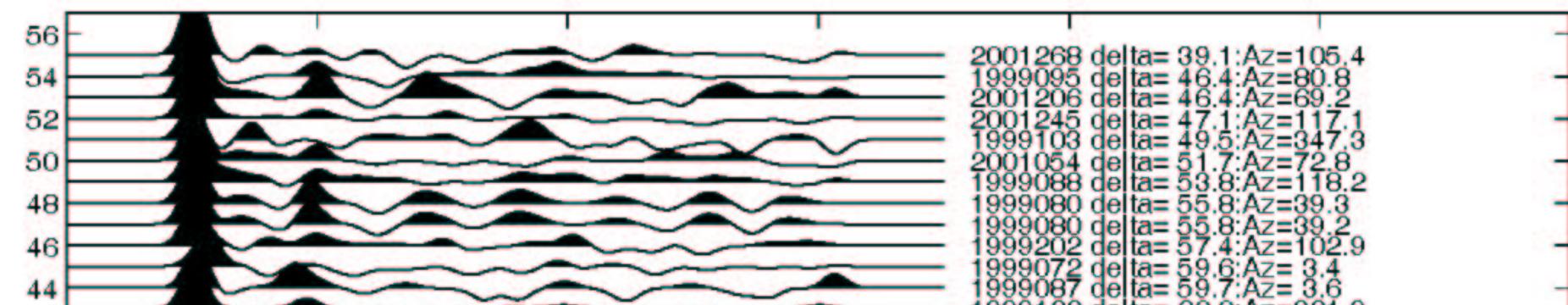
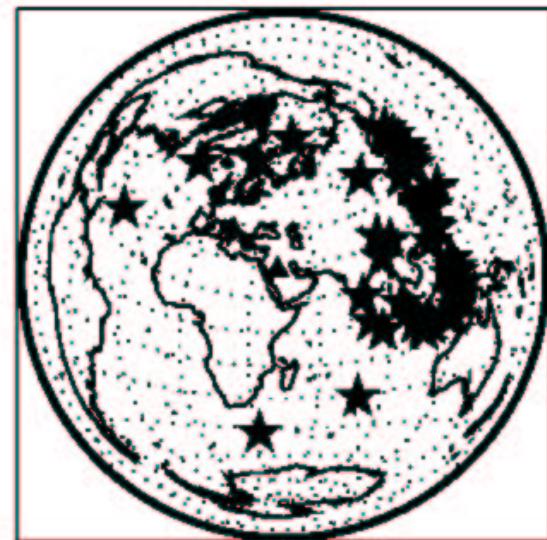
- > consider plane P wave incident from below
- > receiver side scattering includes forward and back scattering, P and S
- > legs ending in P and S isolated through modal decomposition



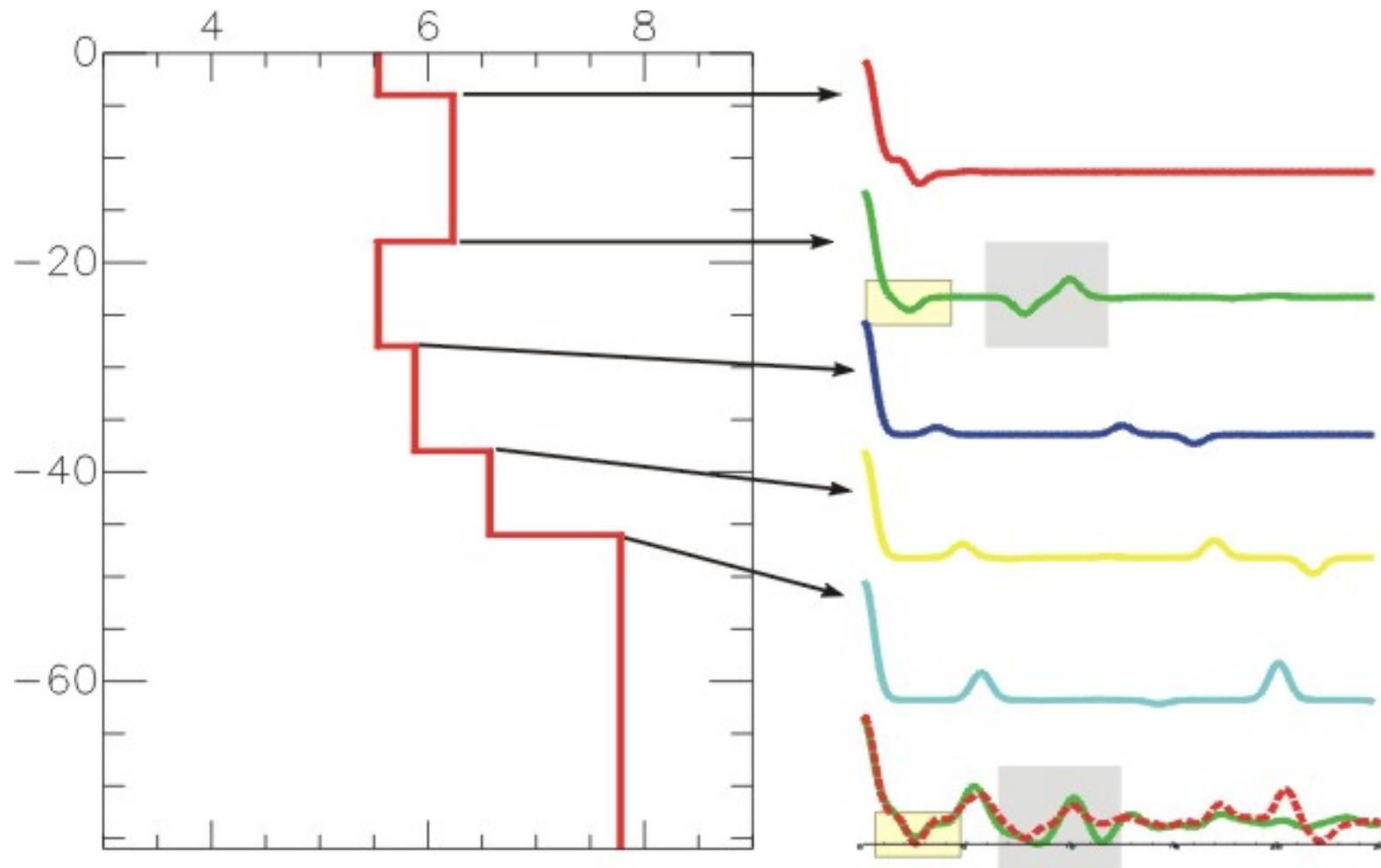
Slant Stacking

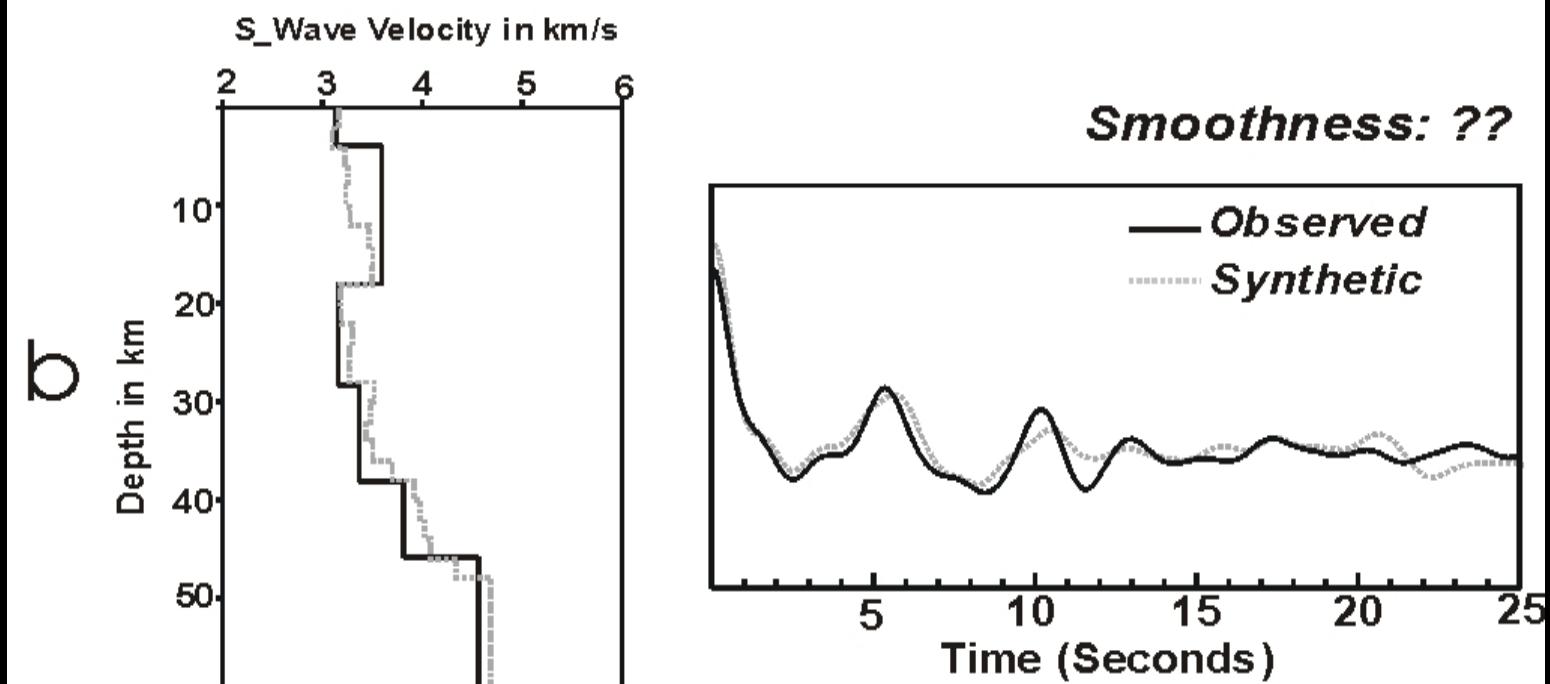
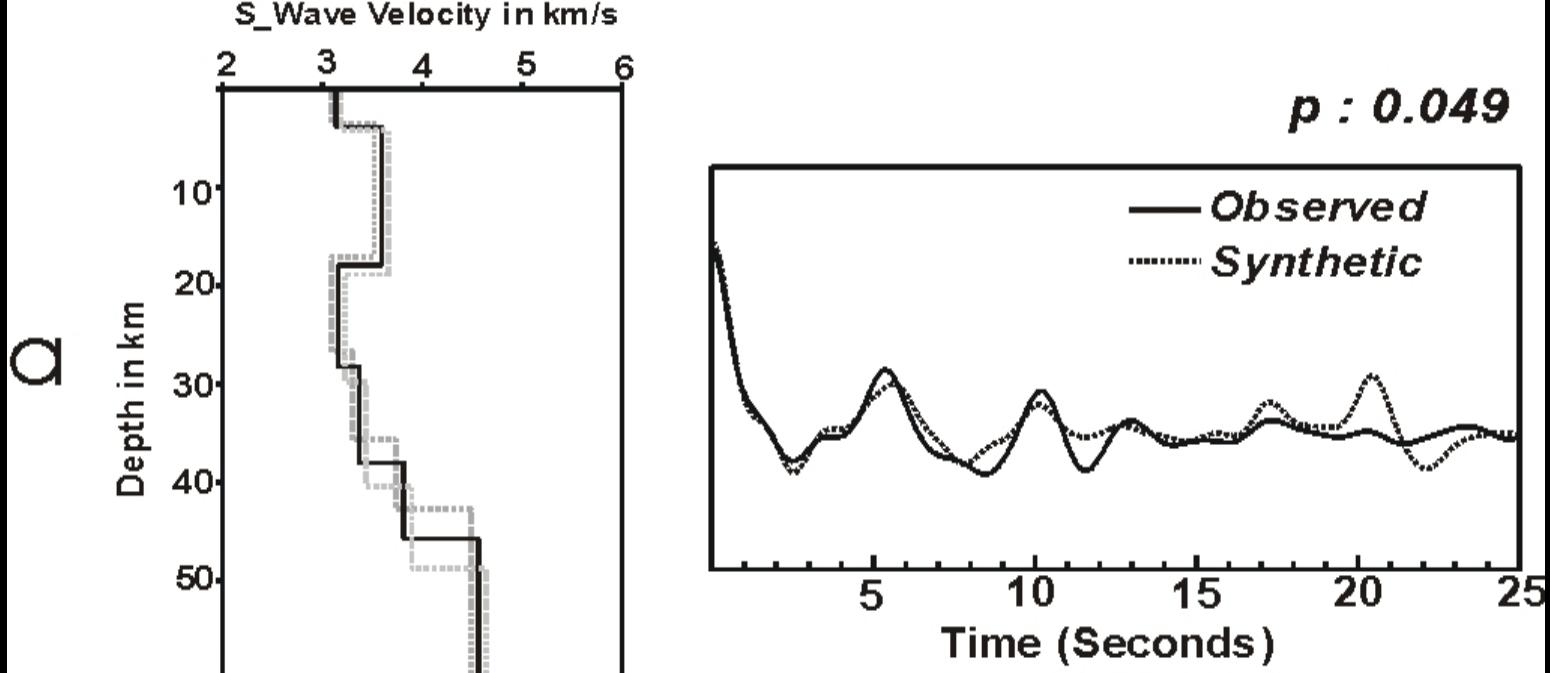


RF map for HILS for a total of (55) EQs



Synthetic Receiver Functions:

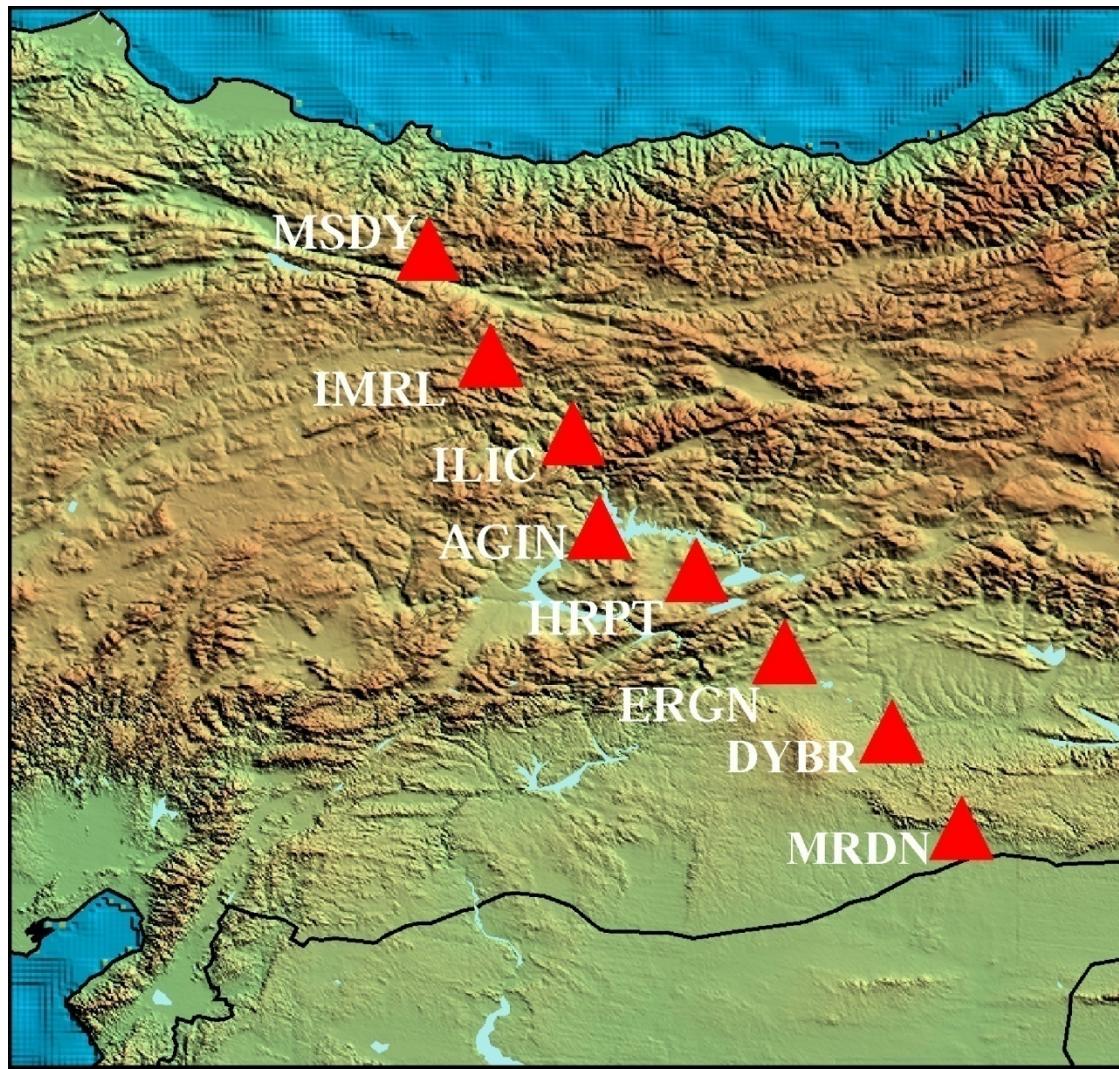




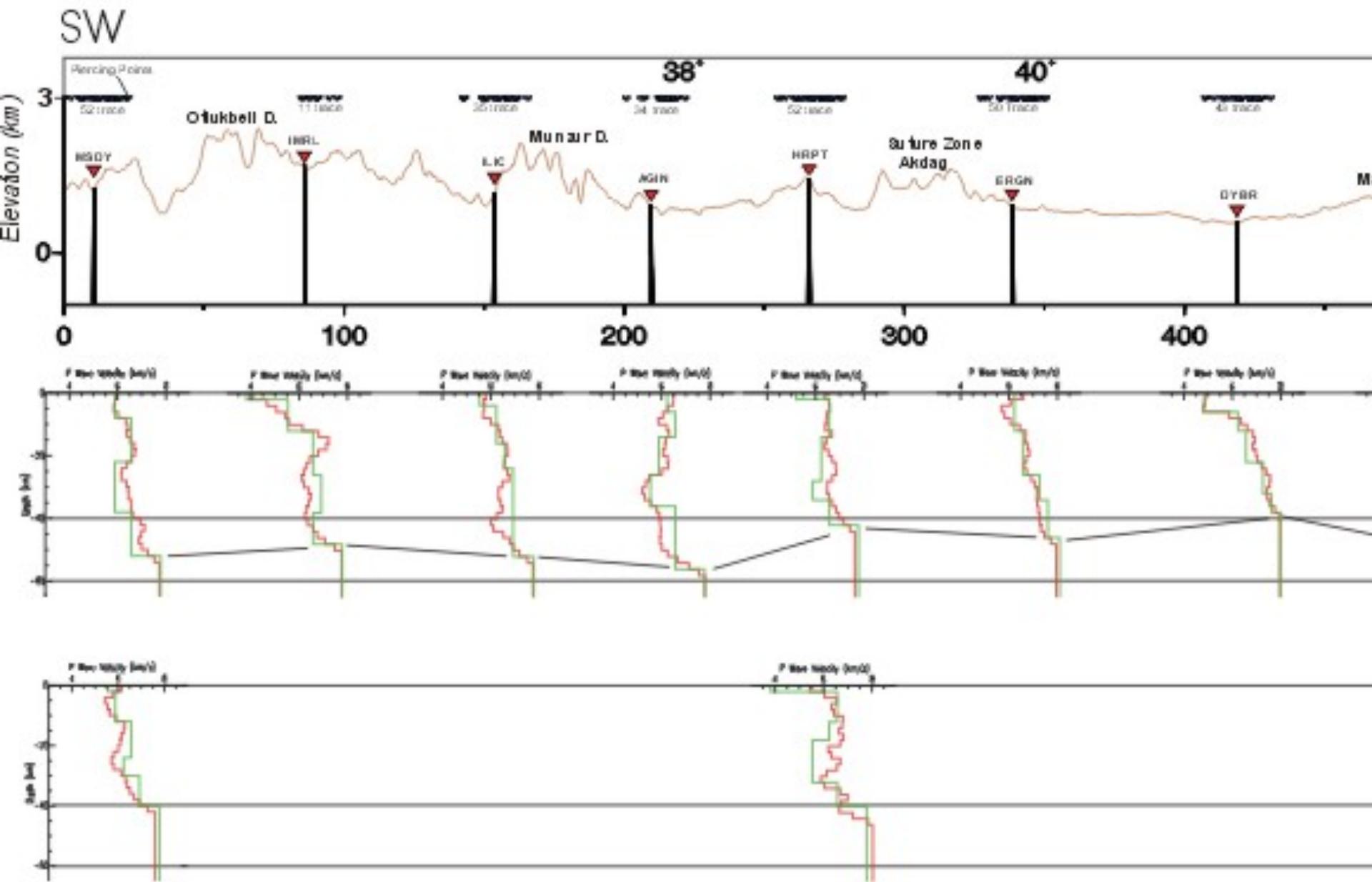
LEAST-SQUARES INVERSION

- > receiver function inversion cast in standard inverse theory framework
- > less expensive than MC/DS methods and makes less stringent demands on data than inverse scattering methods
- > data insufficiency compensated for by regularization (e.g. damping)
- > like MC/DS methods LS involves model matching so there is no formal requirement that data are delivered as Green's functions (e.g. receiver function is adequate); only a forward modelling engine is strictly required

Receiver Functions

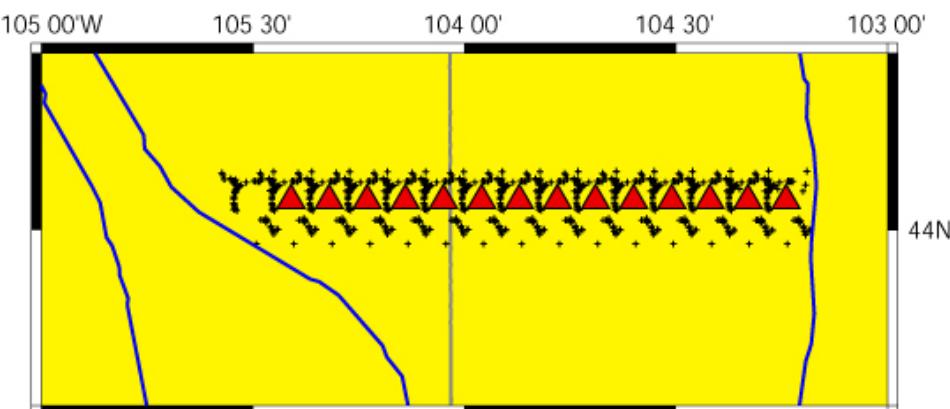


WESTERN PR

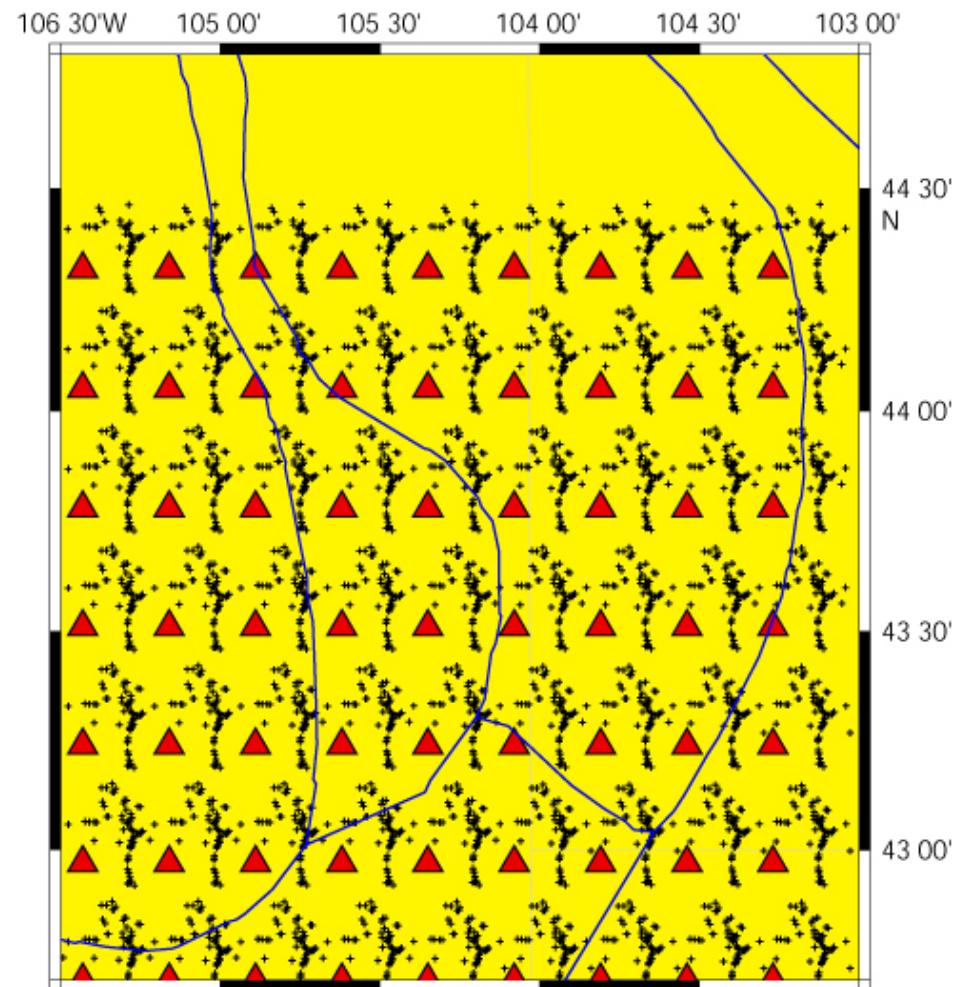


Receiver Function Footprints

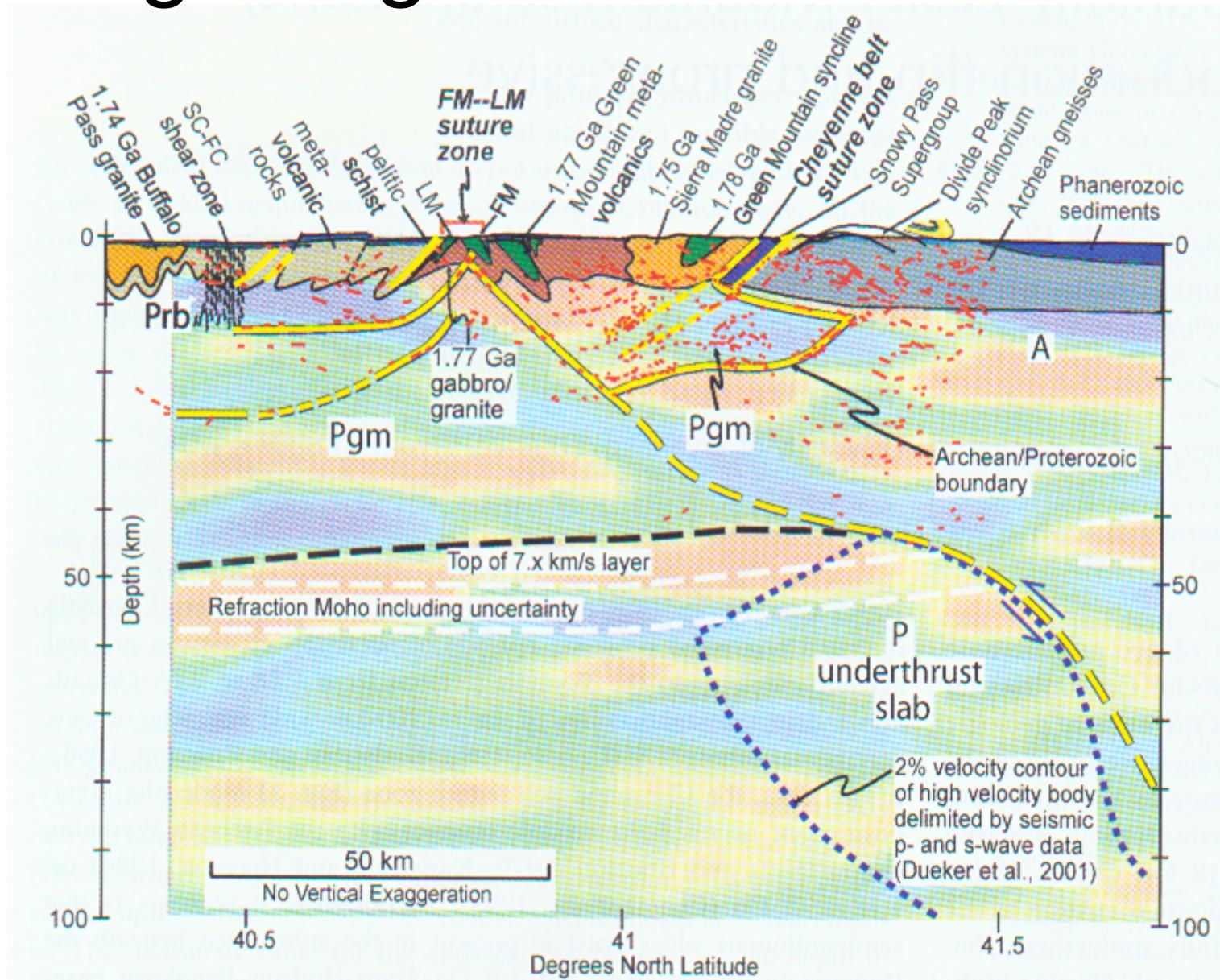
15 km Interface



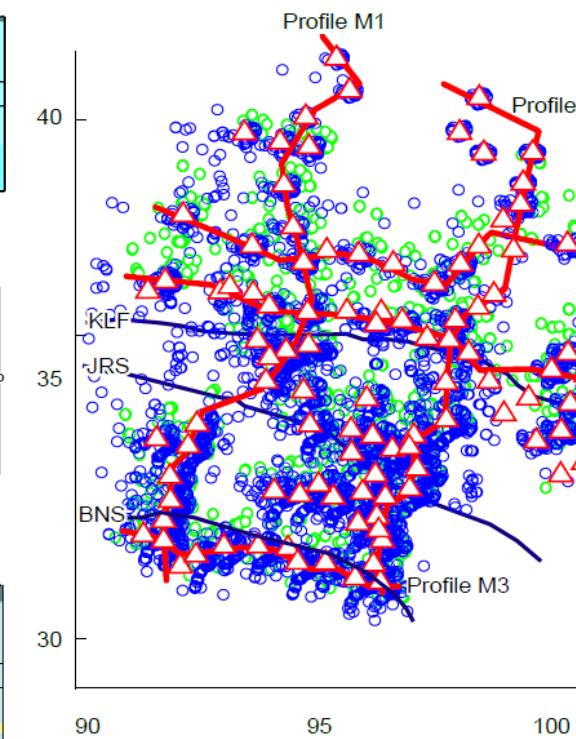
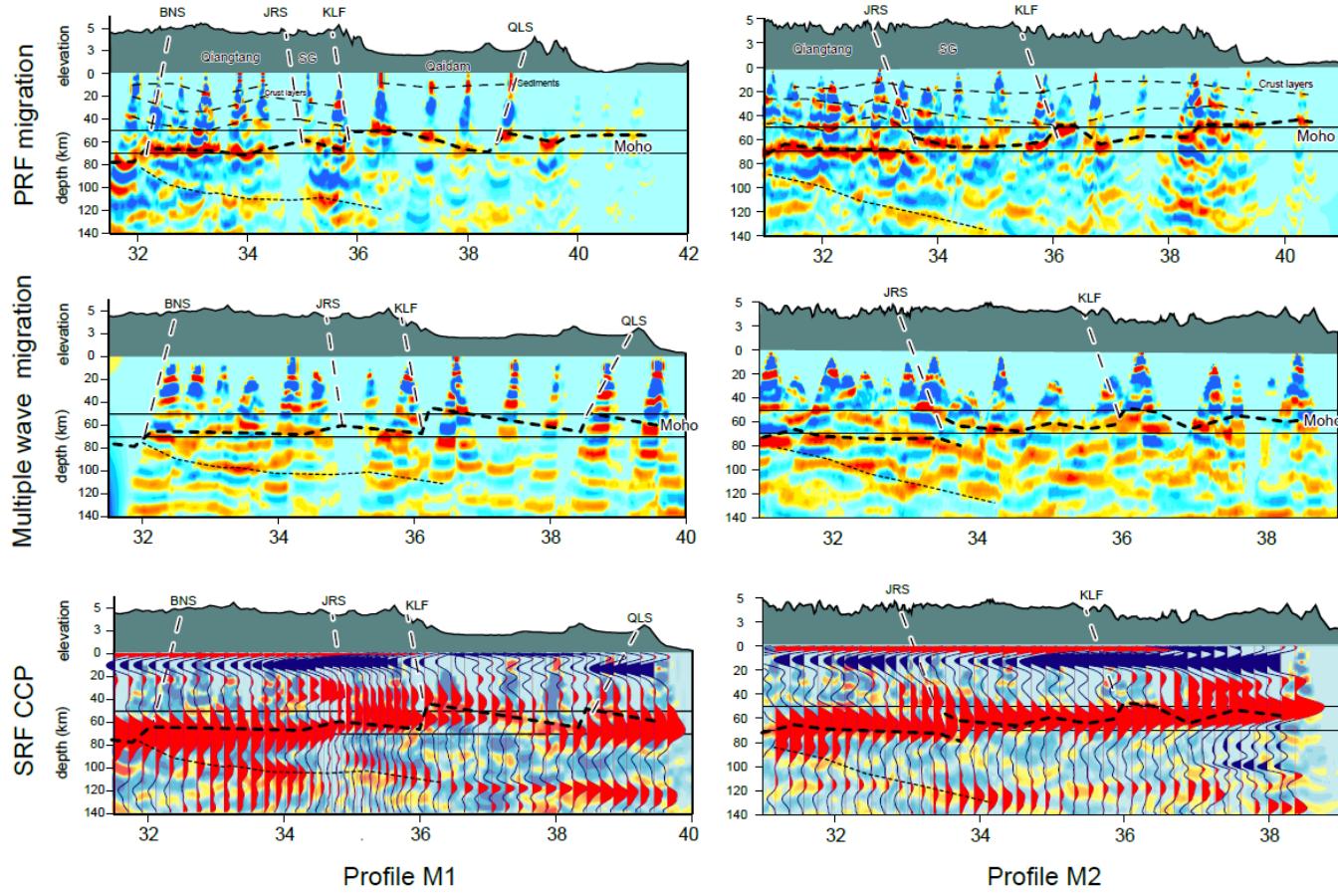
40 km Interface



Migrating Receiver Functions

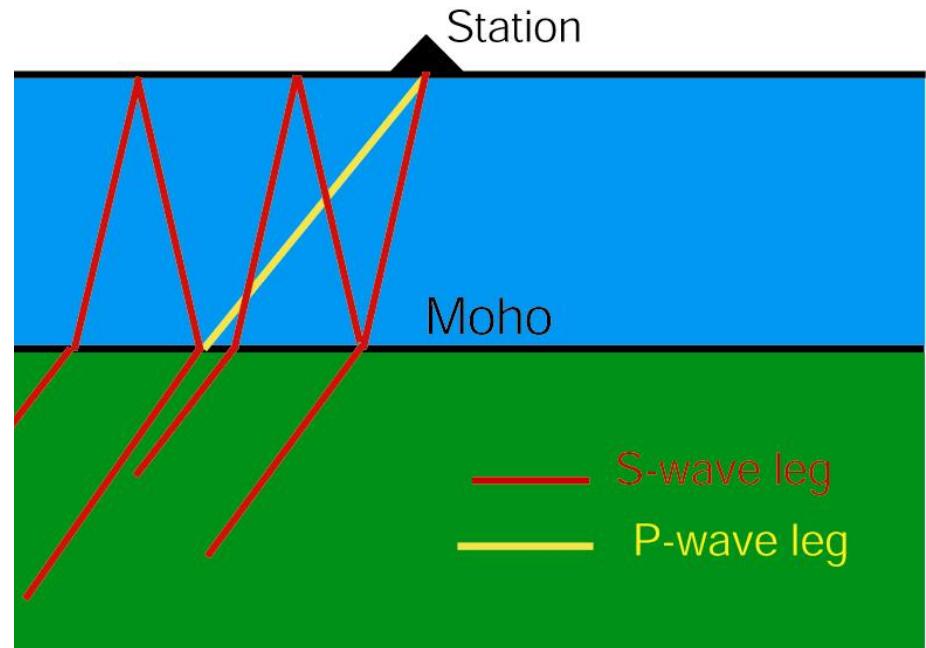
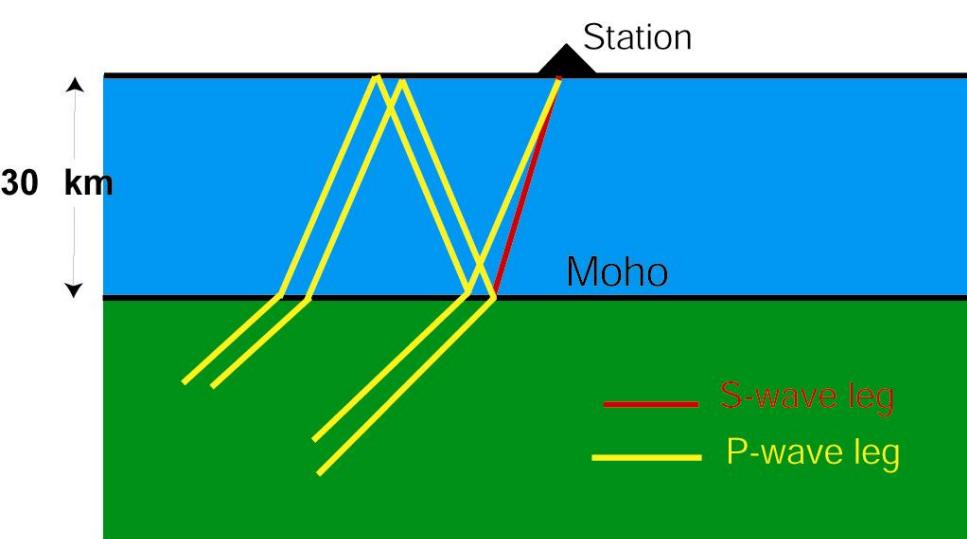


ASCENT P-wave Receiver Functions

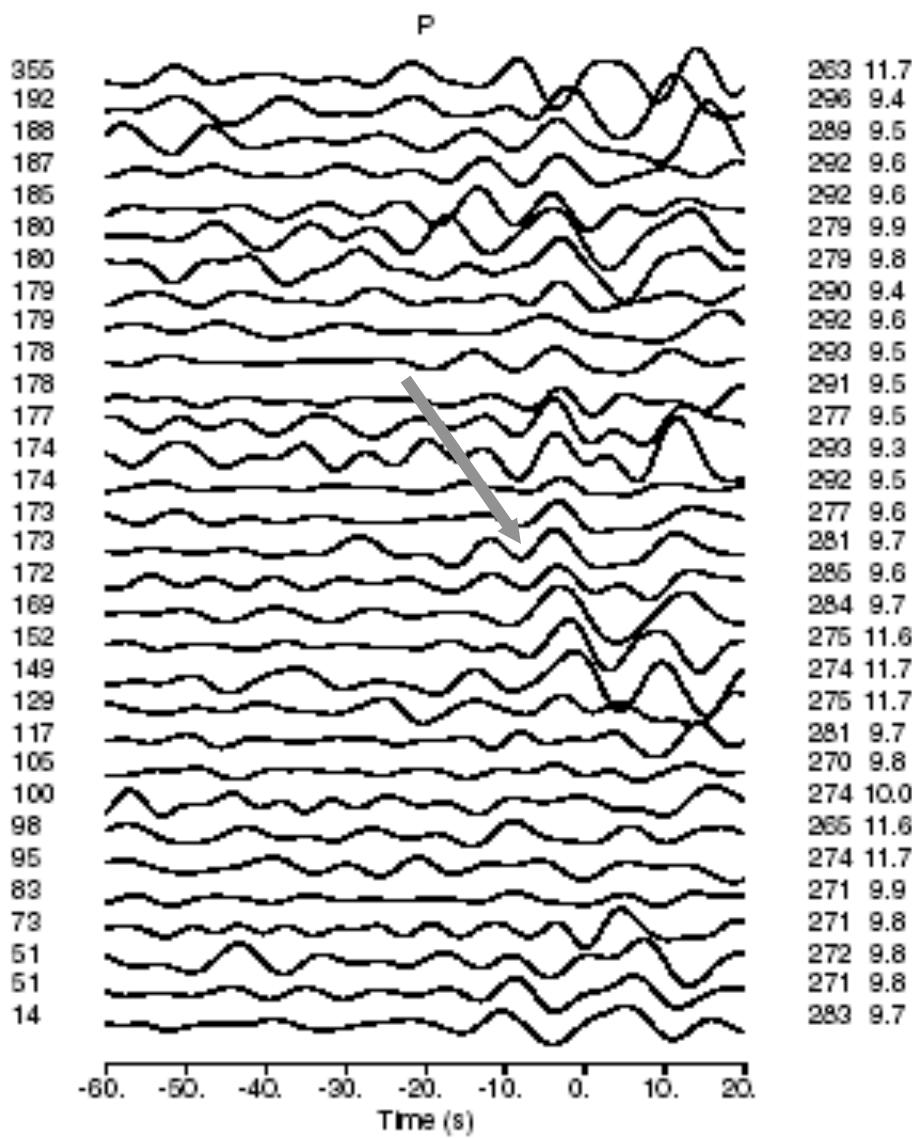
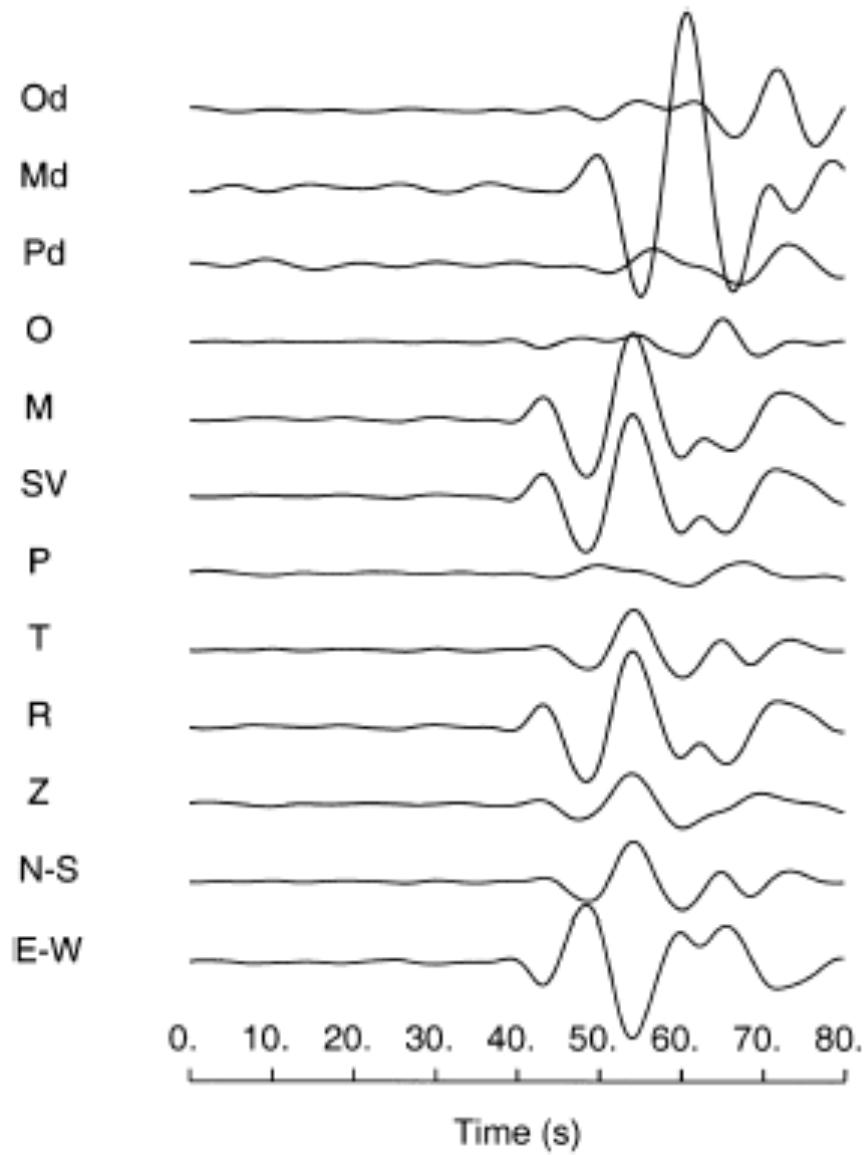


a

S wave Receiver Functions

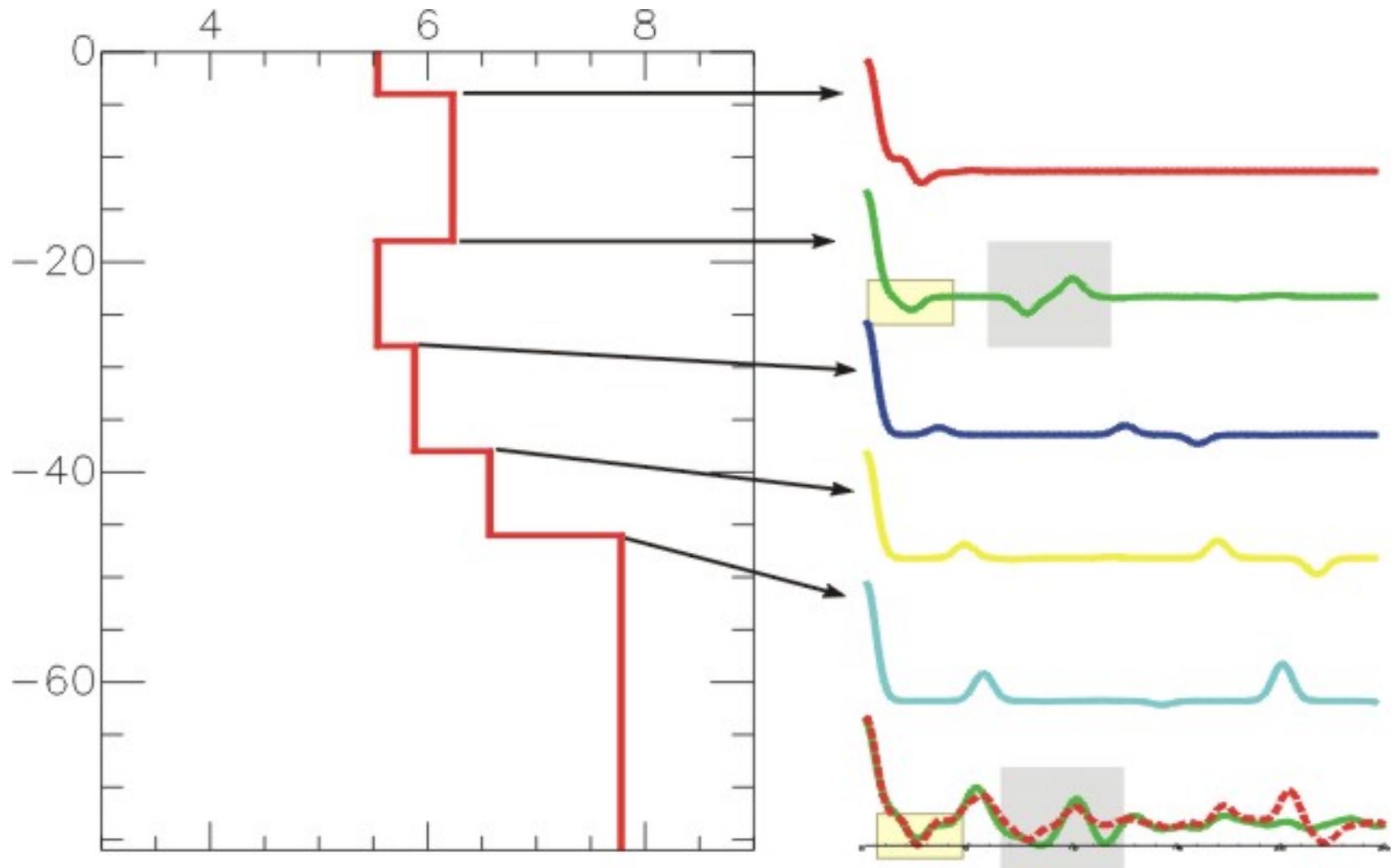


S Receiver Functions

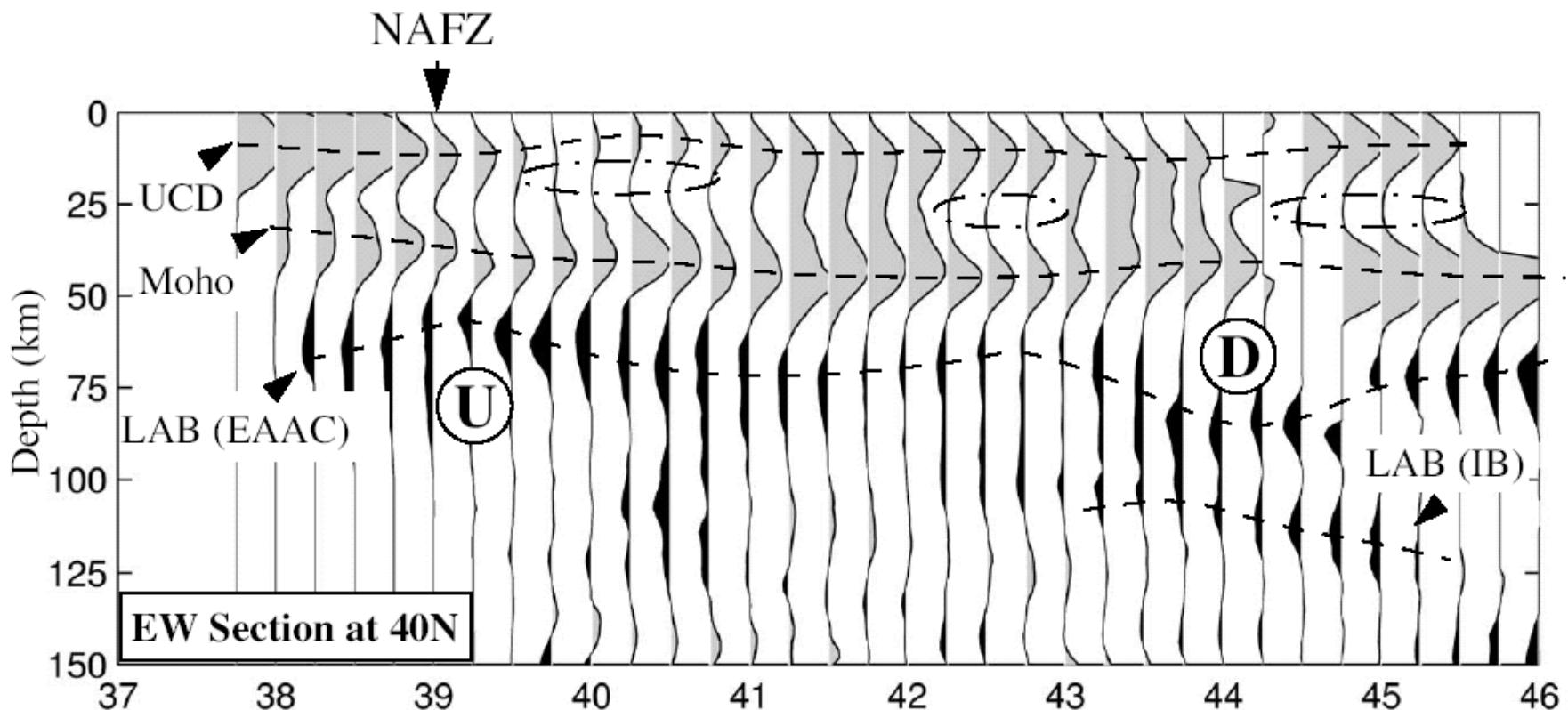


From Farra and Vinnik, 2000

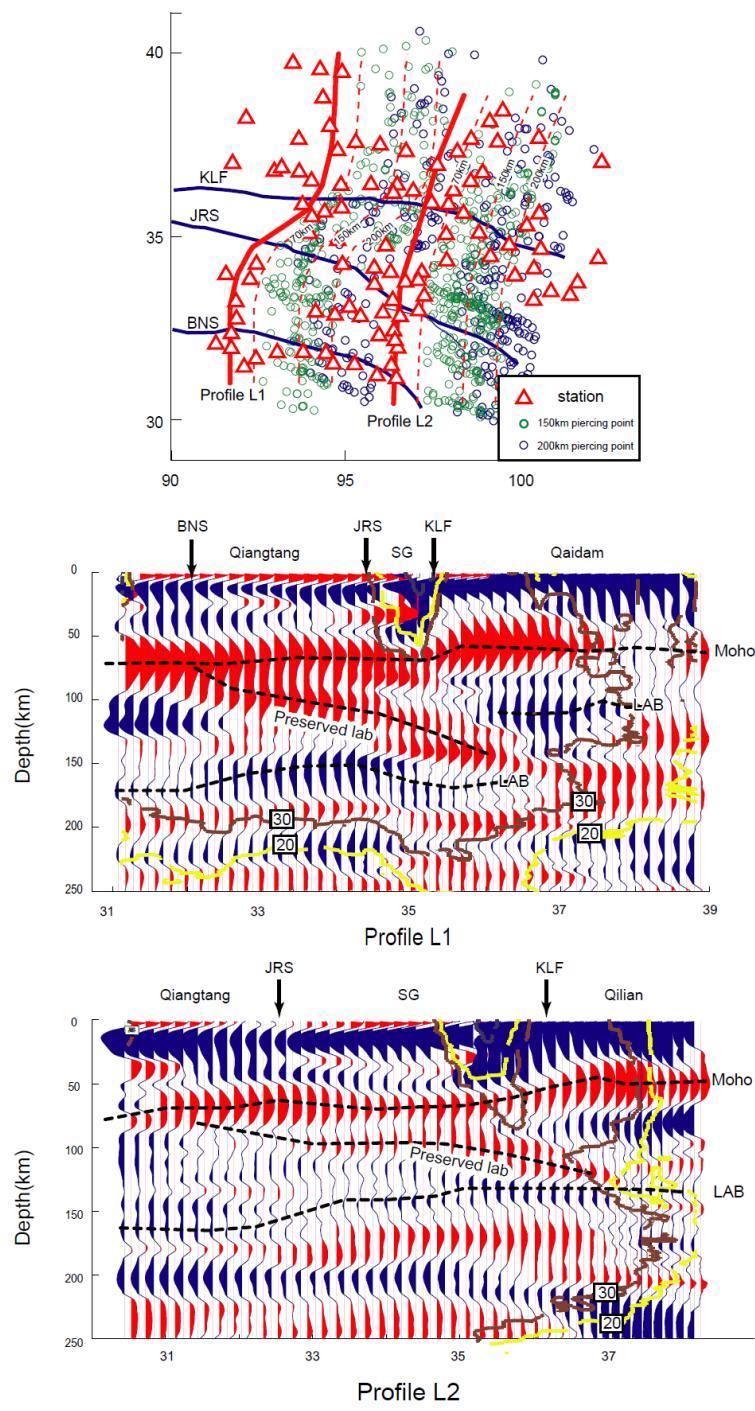
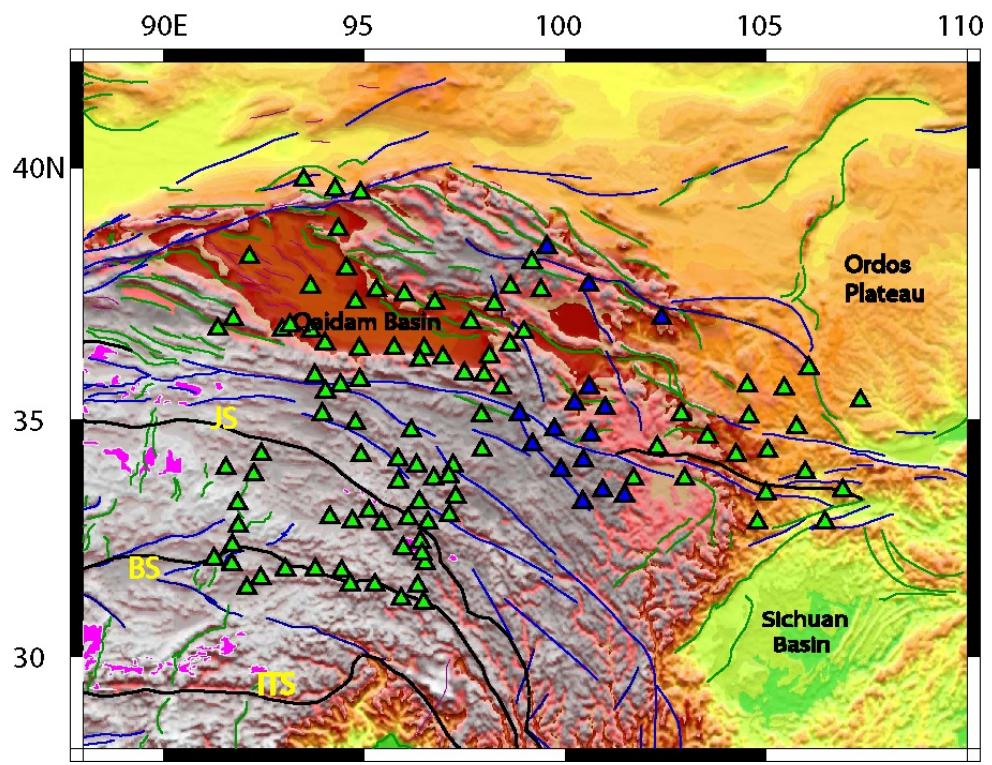
Comparisons with P Receiver Functions



Migrated S-wave Receiver Functions

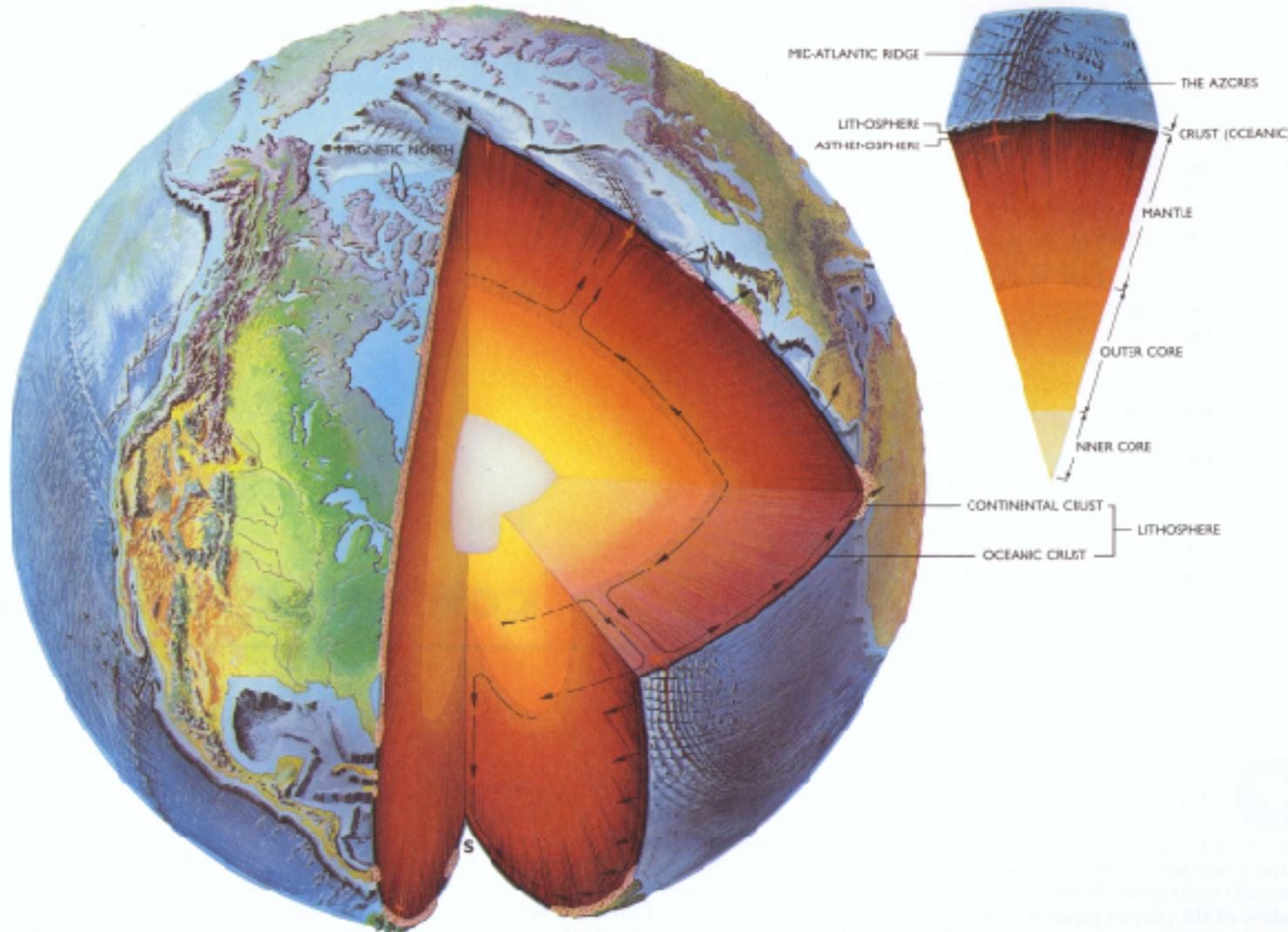


Angus et al., 2005



For Next Time:

Mike Pasyanos: Surface Wave Dispersion

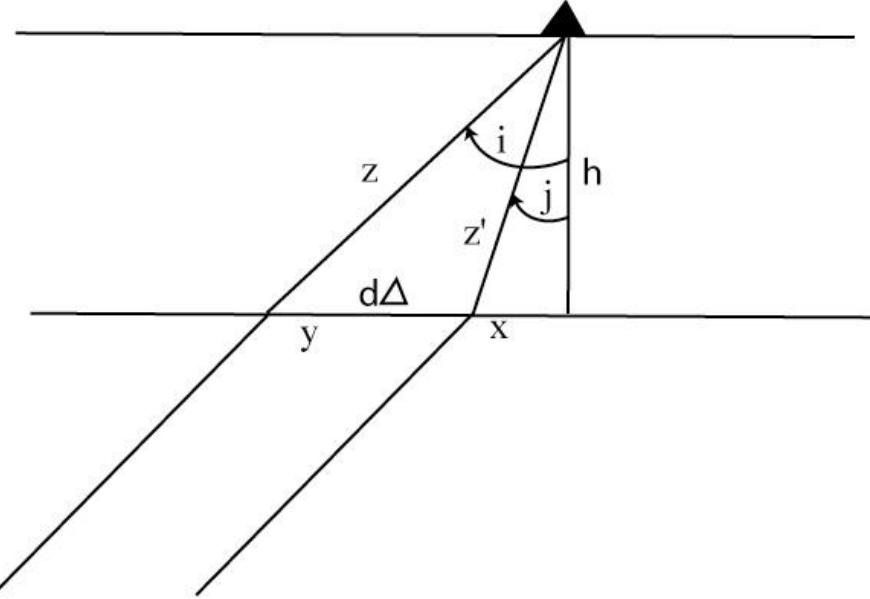


SIMULTANEOUS DECONVOLUTION

- > when large numbers of seismograms representing a single receiver/Green's function are available, perform simultaneous, least-squares deconvolution

$$R(t) = \mathcal{F}^{-1} \left\{ \frac{\sum_i P_i^*(\omega) S_i(\omega)}{\sum_i P_i^*(\omega) P_i(\omega) + \delta} \right\}$$

- > advantageous due to fact that smaller sum of spectra in denominator reduce likelihood of spectral zeros allowing for smaller values of water level parameter δ to be used
- > Gurrola et al, 1995, GJI, 120, 537-543



$$t_{PS-P} = \frac{h(V_p \sqrt{1 - p^2 V_s^2} - V_s \sqrt{1 - p^2 V_p^2})}{V_s V_p \sqrt{1 - p^2 V_p^2} \sqrt{1 - p^2 V_s^2}} - d\Delta p$$

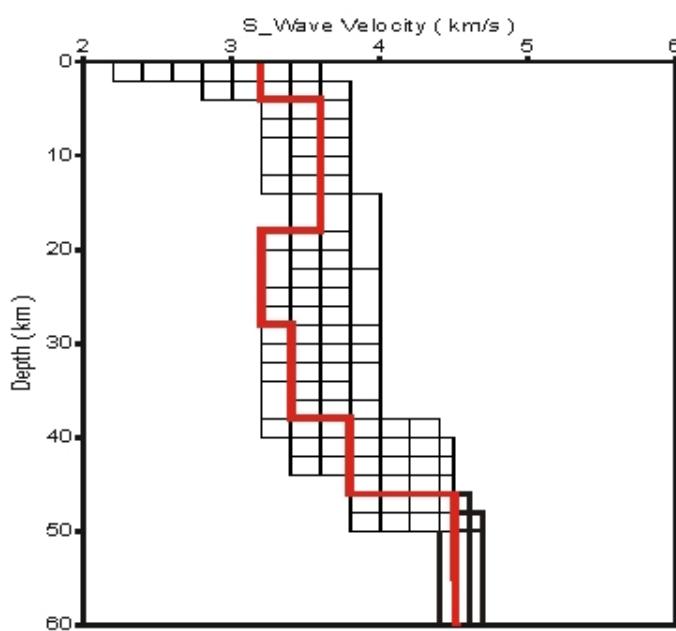
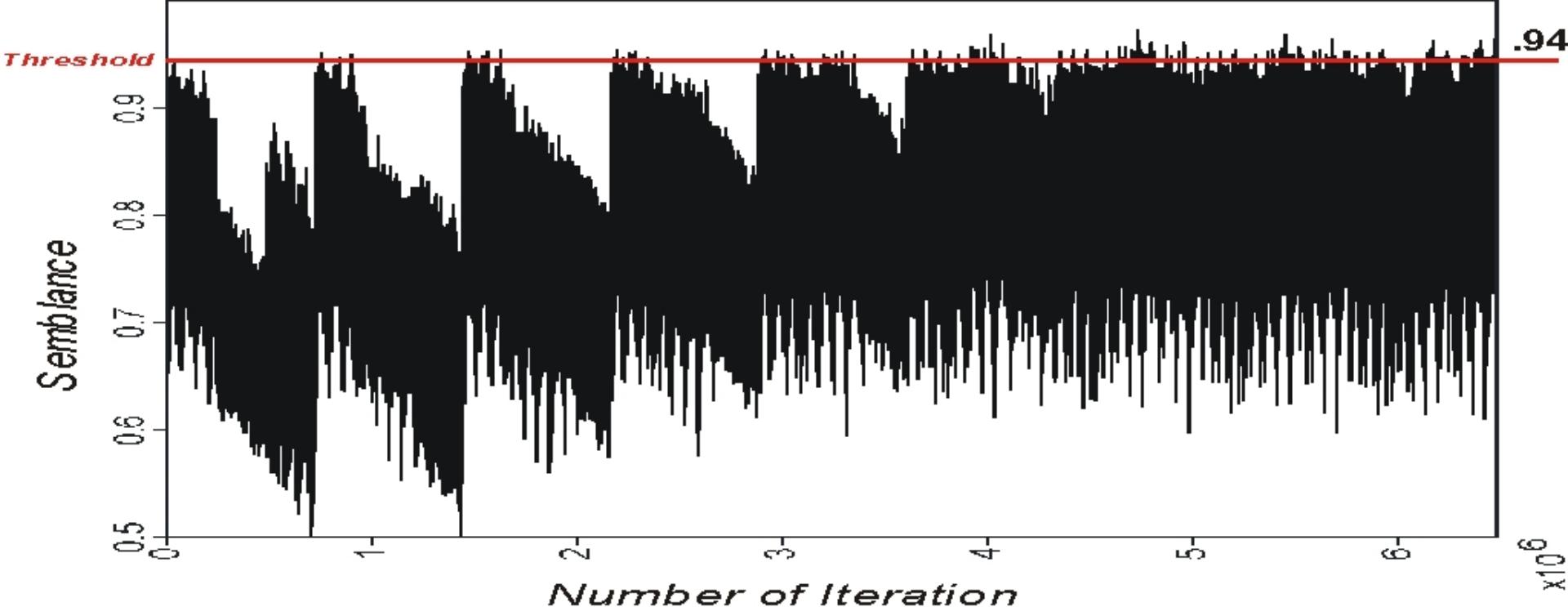
$$= \frac{h(V_p \sqrt{1 - p^2 V_s^2} - V_s \sqrt{1 - p^2 V_p^2})}{V_s V_p \sqrt{1 - p^2 V_p^2} \sqrt{1 - p^2 V_s^2}} + p(y - x)$$

$$x = h \tan j = \frac{h V_s p}{\sqrt{1 - p^2 V_s^2}}, y = h \tan i = \frac{h V_p p}{\sqrt{1 - p^2 V_p^2}}$$

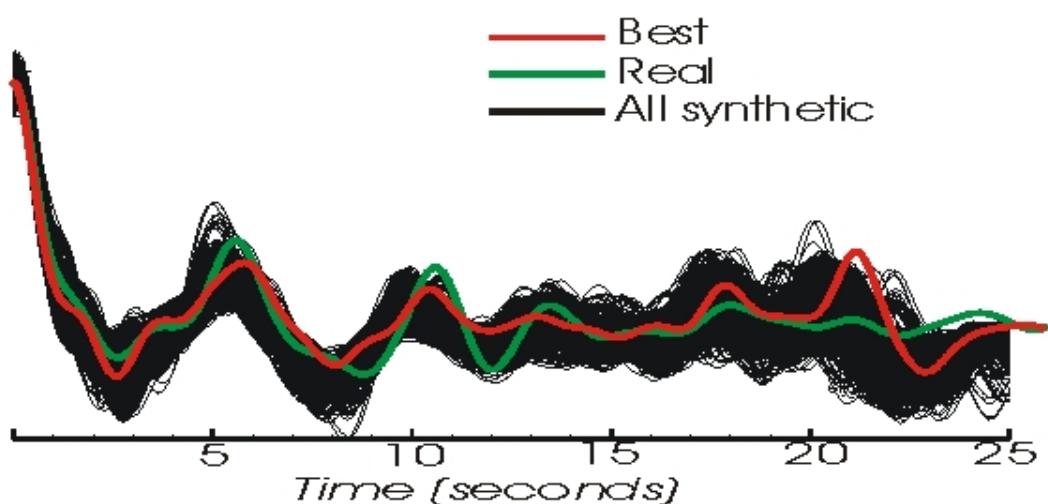
$$\therefore t_{PS-P} = \frac{h(V_p \sqrt{1 - p^2 V_s^2} - V_s \sqrt{1 - p^2 V_p^2})}{V_s V_p \sqrt{1 - p^2 V_p^2} \sqrt{1 - p^2 V_s^2}} + h p^2 \frac{V_p \sqrt{1 - p^2 V_s^2} - V_s \sqrt{1 - p^2 V_p^2}}{\sqrt{1 - p^2 V_s^2} \sqrt{1 - p^2 V_p^2}}$$

$$= h \left[\frac{(V_p \sqrt{1 - p^2 V_s^2} - V_s \sqrt{1 - p^2 V_p^2}) + p^2 V_p^2 V_s \sqrt{1 - p^2 V_s^2} - p^2 V_s^2 V_p \sqrt{1 - p^2 V_p^2}}{V_s V_p \sqrt{1 - p^2 V_s^2} \sqrt{1 - p^2 V_p^2}} \right]$$

$$= h \left[\frac{V_p \sqrt{1 - p^2 V_s^2} (1 - p^2 V_s^2) - V_s \sqrt{1 - p^2 V_p^2} (1 - p^2 V_p^2)}{V_s V_p \sqrt{1 - p^2 V_s^2} \sqrt{1 - p^2 V_p^2}} \right] = h \left[\sqrt{V_s^{-2} - p^2} - \sqrt{V_p^{-2} - p^2} \right]$$

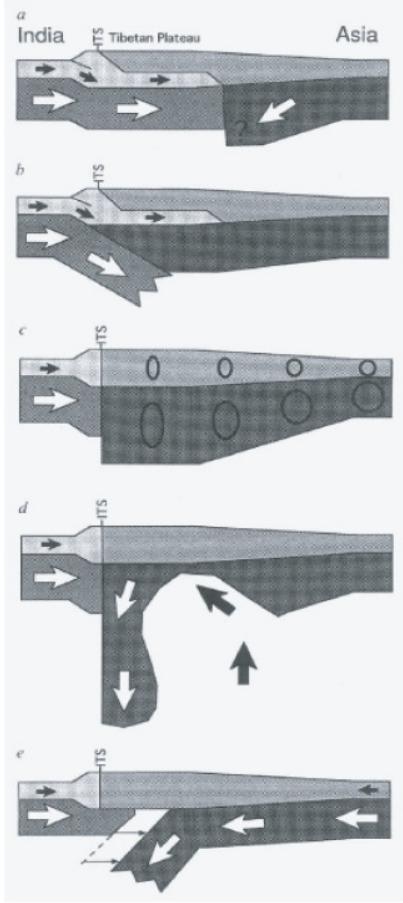


1069 Model

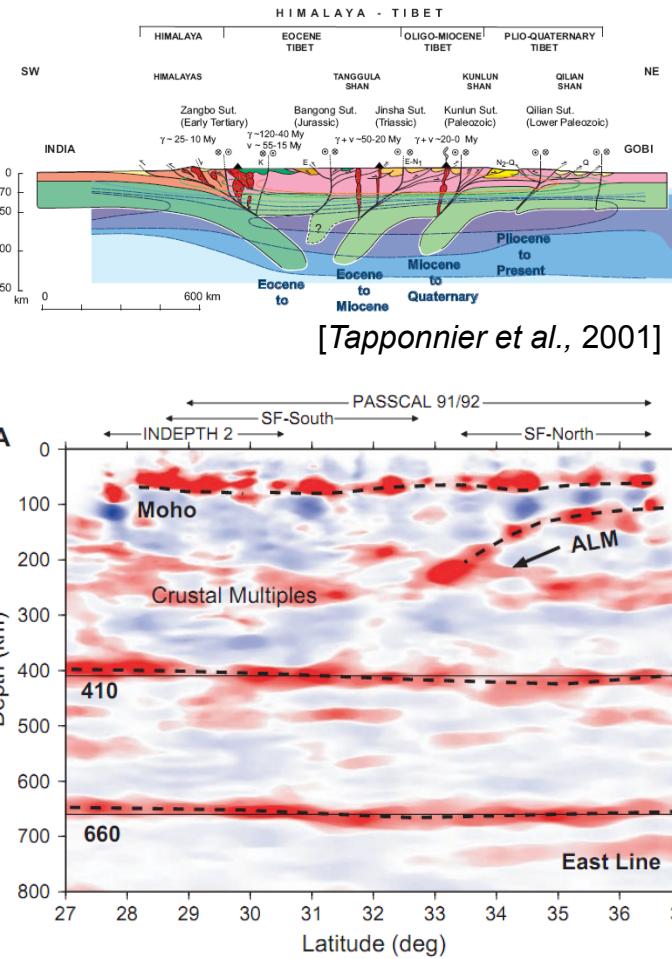


Semblance Fit: 0.96

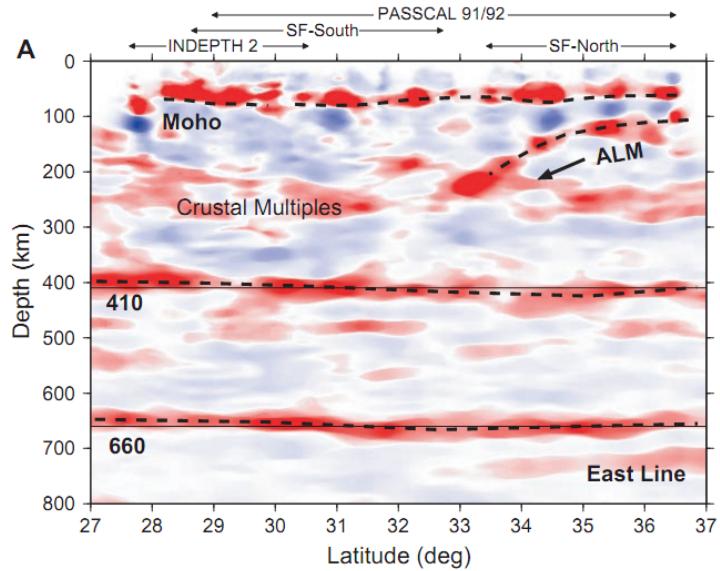
Background: Geodynamic models of Tibetan plateau



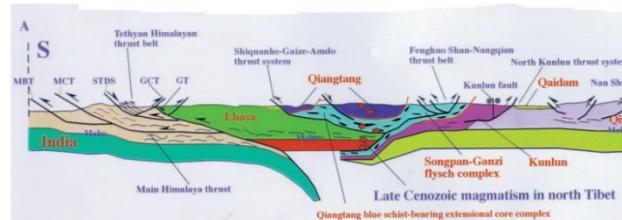
[Willett and Beaumont, 1994]



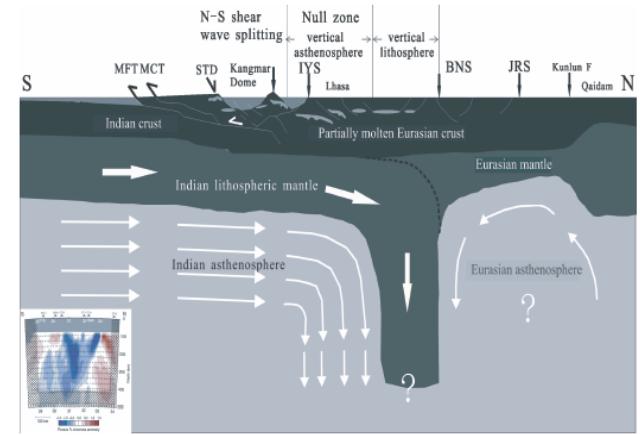
[Tappognier et al., 2001]



[Kind et al., 2002]



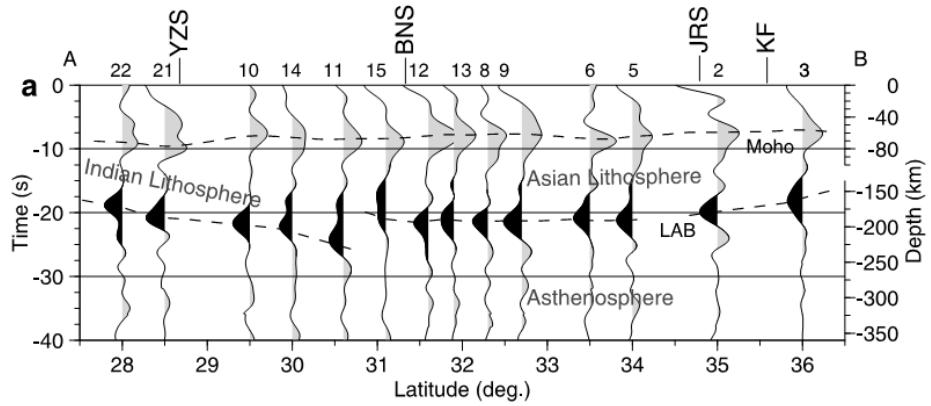
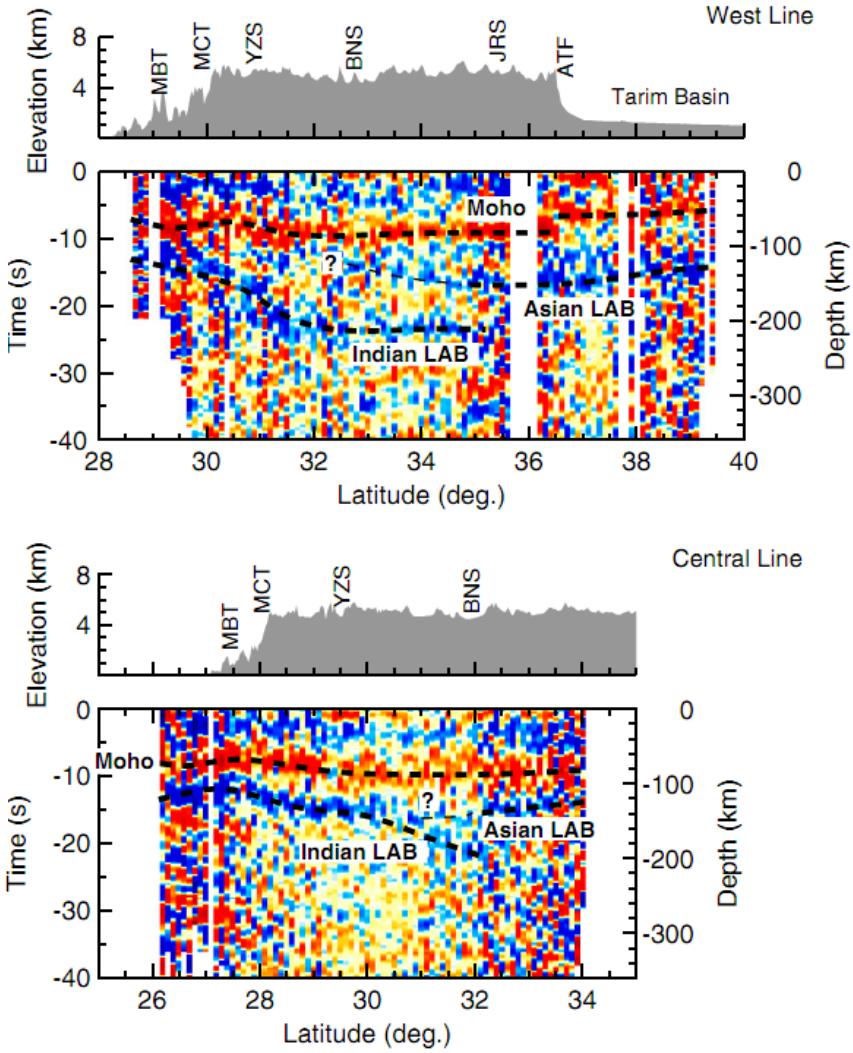
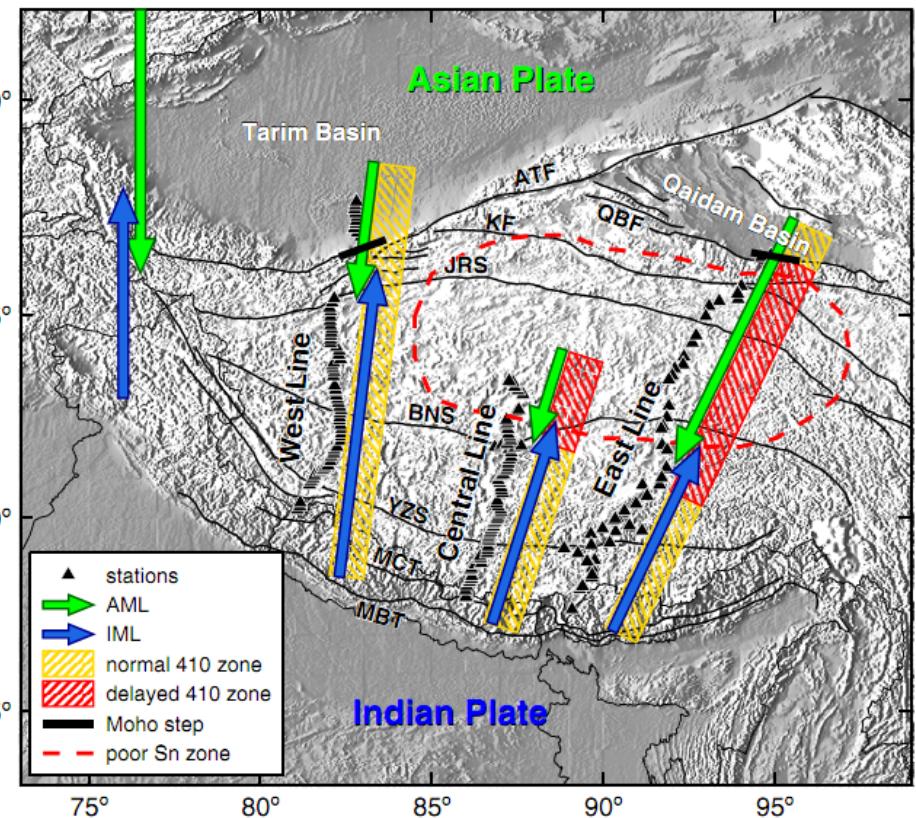
[Yin and Harrison, 2000]



[Fu et al., 2008]

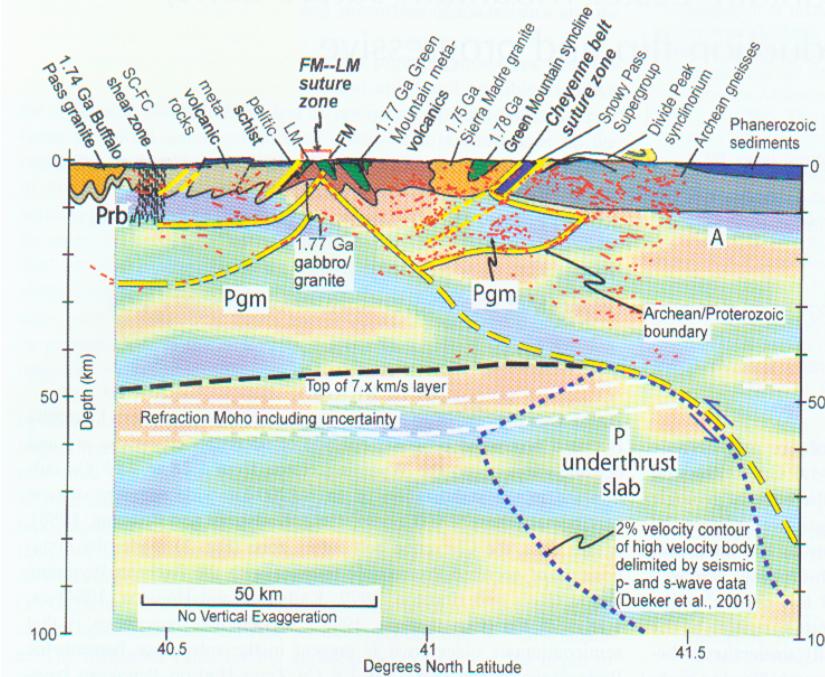
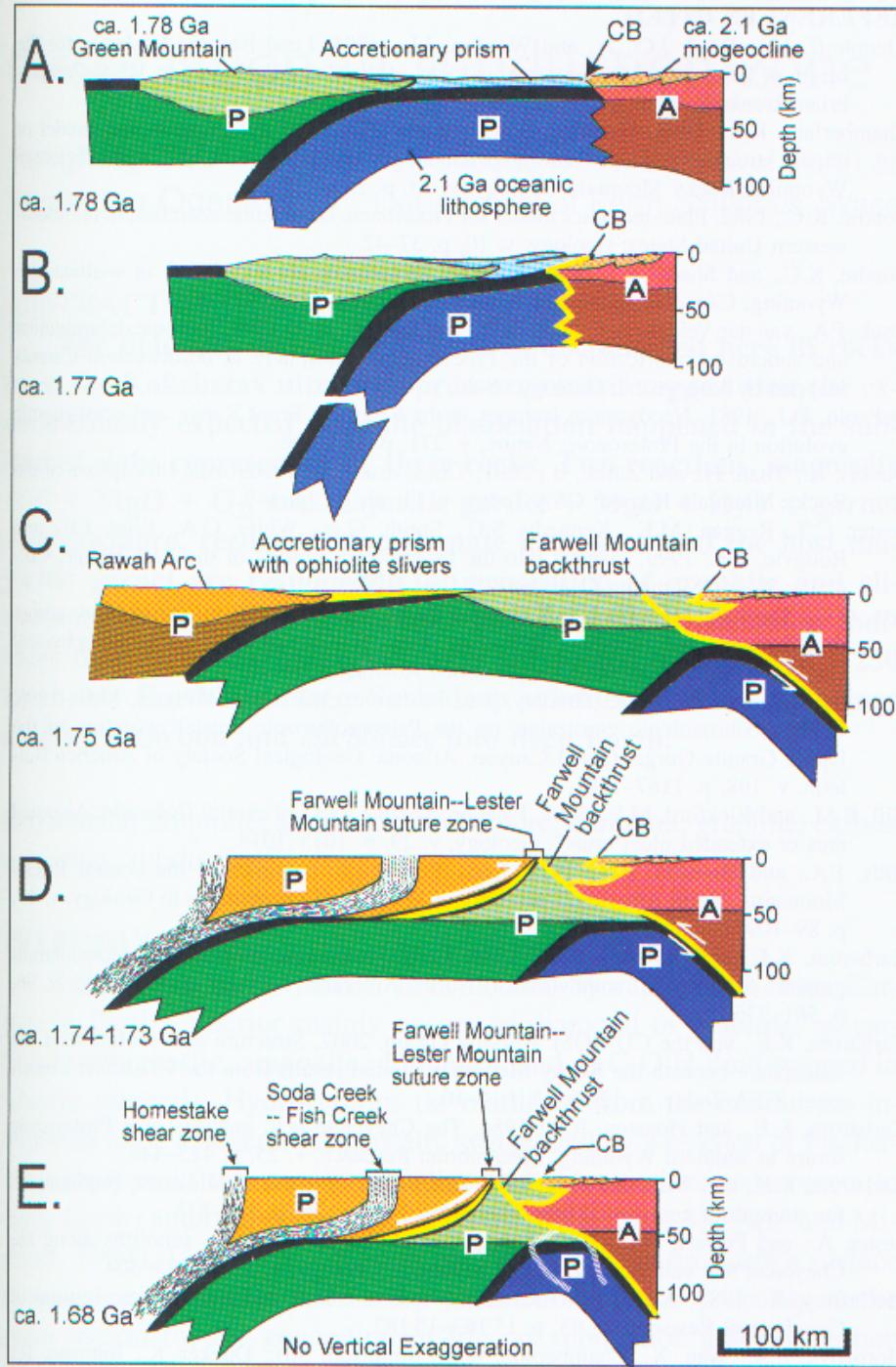
Unknown deep structure

Mainly north-south profiles



[Kumar et al.,

[Zhao et al., 2010]



Migrated S-wave Receiver Functions

